

Proton & Neutron Polarizabilities with Compton Scattering from Low-Mass Nuclear Targets at the High Intensity Gamma-ray Source (HIGS)

Kent Leung (Nuclear Photonics 2025, Darmstadt, Germany)

Assist. Prof., Physics & Astronomy Dept., Montclair, New Jersey



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View 20 km in this direction



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Compton @ HIGS

Nuclear Compton Scattering @ the High-Intensity Gamma-ray Source

Duke University: J. Zhou, E. Mancil, F. Friesen, H. Gao, D. Godagama, C. R. Howell, S. Jia, S. Mikhailov, Y. K. Wu, B. Yu, C. Martin

North Carolina Central University: M. Ahmed, B. Crowe, D. Markoff

George Washington University: E. Downie, G. Feldman, M. Lewis,

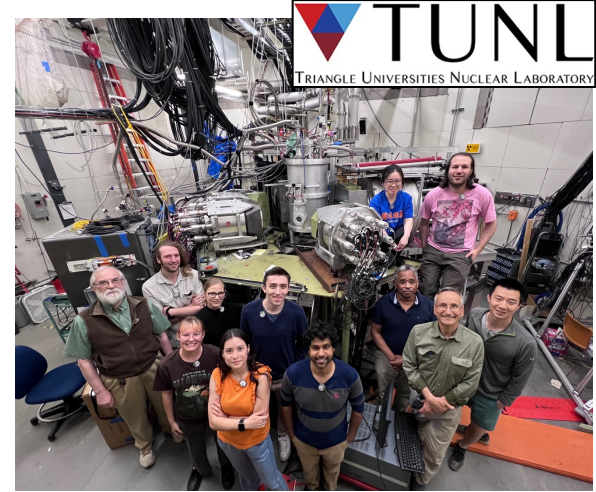
Mount Allison University: D. Hornidge

University North Carolina at Chapel Hill: H. Karwowski

University of Kentucky: M. Kovash

Montclair State University: K. Leung, S. Estupinan Jimenez

University of Saskatchewan: R. Pywell

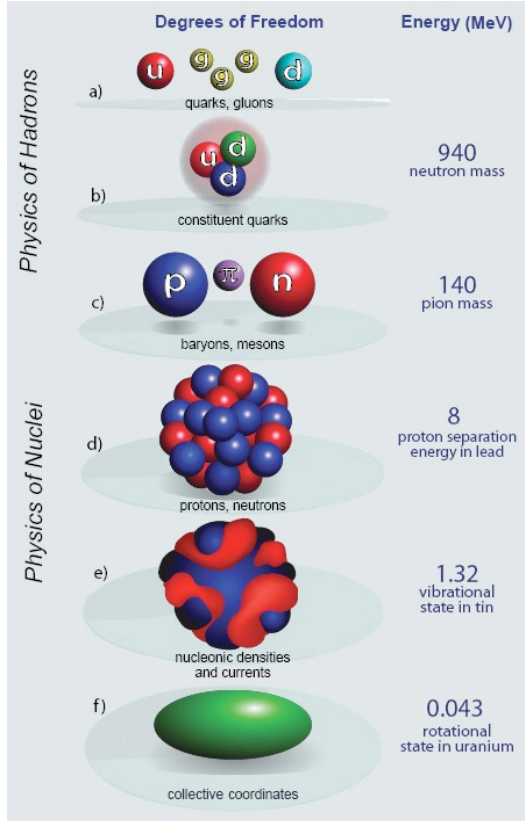


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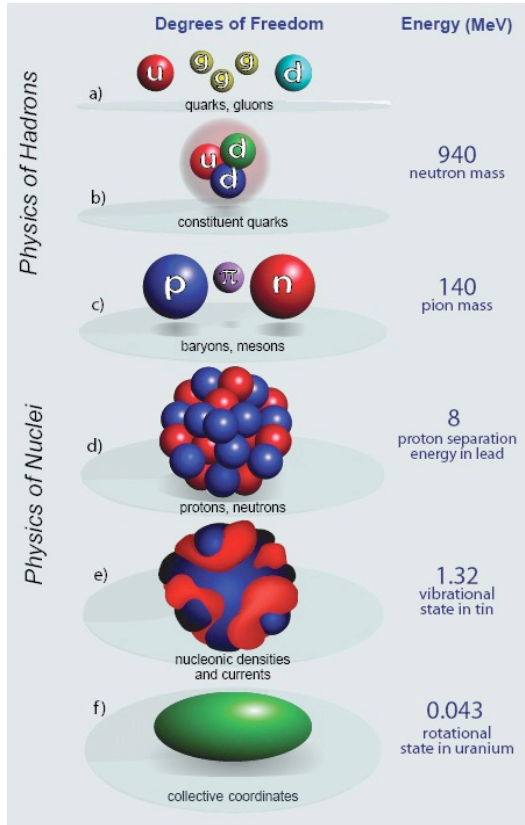
Low-energy QCD with nucleon polarizabilities

- QCD Asymptotic Freedom. Gross, Wilczek & Politzer. Physics Nobel 2004. Non-perturbative.

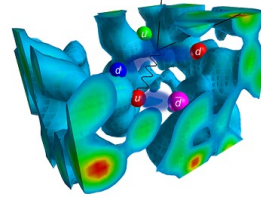


Low-energy QCD with nucleon polarizabilities

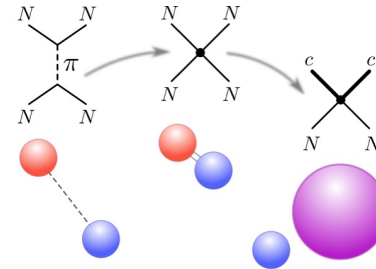
- QCD Asymptotic Freedom. Gross, Wilczek & Politzer. Physics Nobel 2004. Non-perturbative.



Lattice QCD

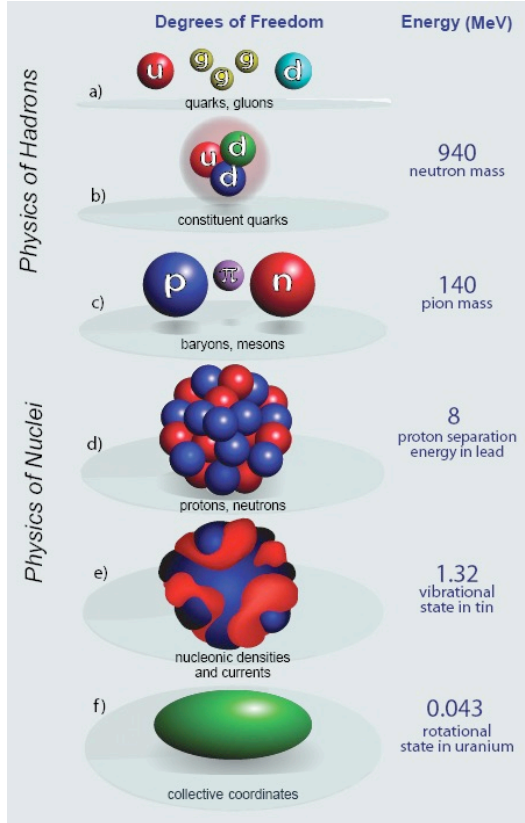


Chiral Effective Field Theory

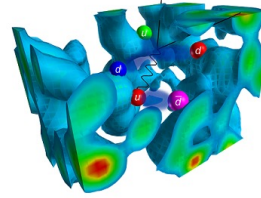


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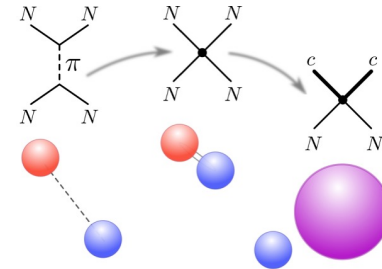
Lattice QCD



Nucleon polarizabilities



Chiral Effective Field Theory



Nucleon polarizabilities

- Fundamental properties probing nucleon structure
- Induced dipole moments from static fields:

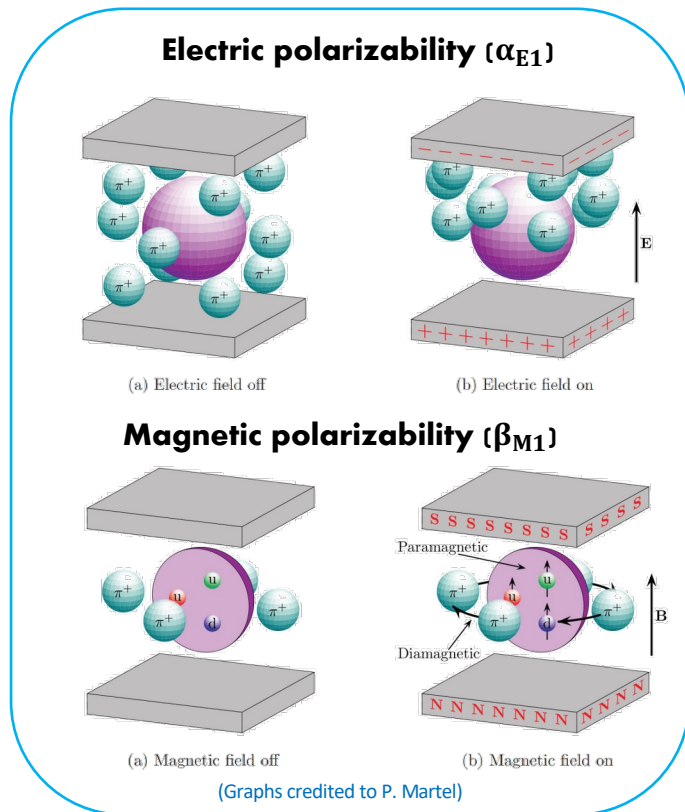
$$\vec{d}_{\text{ind}} = \alpha_{E1} \vec{E}, \quad \vec{m}_{\text{ind}} = \beta_{M1} \vec{H}$$

scalar electromagnetic polarizabilities

Why scalar polarizabilities?

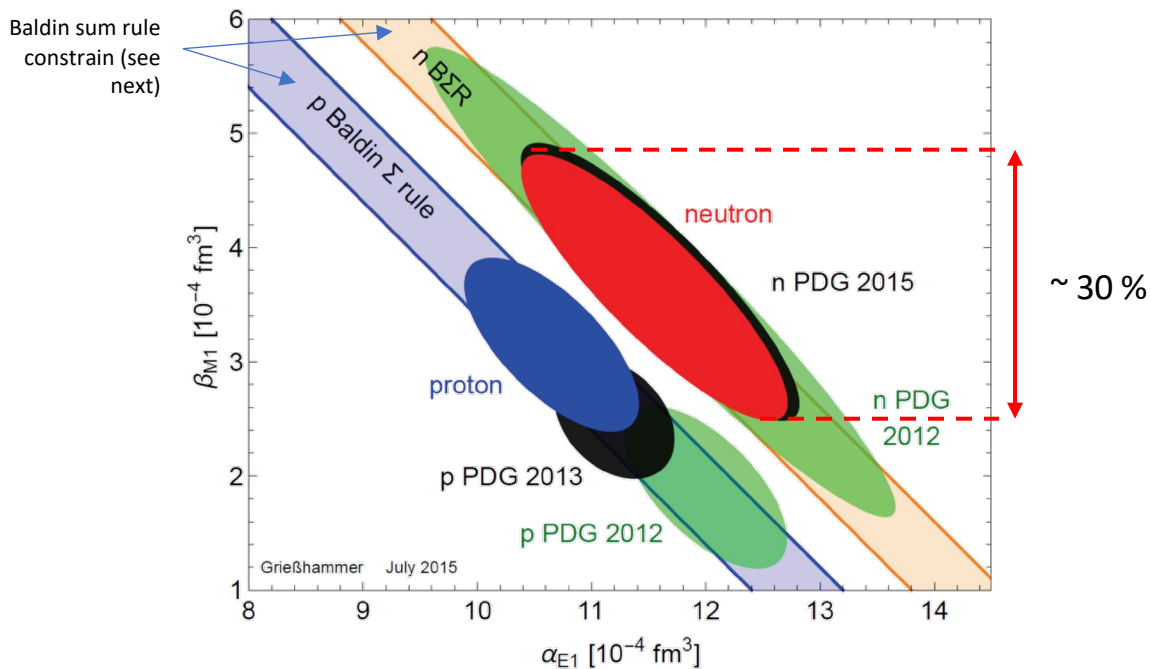
- Low-energy QCD: Bridge between emerging lattice QCD calculations & χ EFT
- Input to Lamb shift in muonic hydrogen for proton radius puzzle
- Electromagnetic contribution to charge symmetry breaking: uncertainty dominated by α_{E1}^{p-n} , β_{M1}^{p-n}

Thanks to Jingyi Zhou



- e.g., the β_M (magnetic polarizability) of a nucleon is an interplay between diamagnetic charged pion currents and paramagnetic Δ resonance

Current knowledge of scalar polarizabilities



Scalar Dipole Polarisabilities: “canonical units” $[10^{-4} \text{ fm}^3]$

$\alpha_{E1} [10^{-4} \text{ fm}^3]$

$\beta_{M1} [10^{-4} \text{ fm}^3]$

proton (Baldin, $N^2\text{LO}$)
McGovern/Phillips/hg EPJA 2013

$10.65 \pm 0.35_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.3_{\text{theory}}$

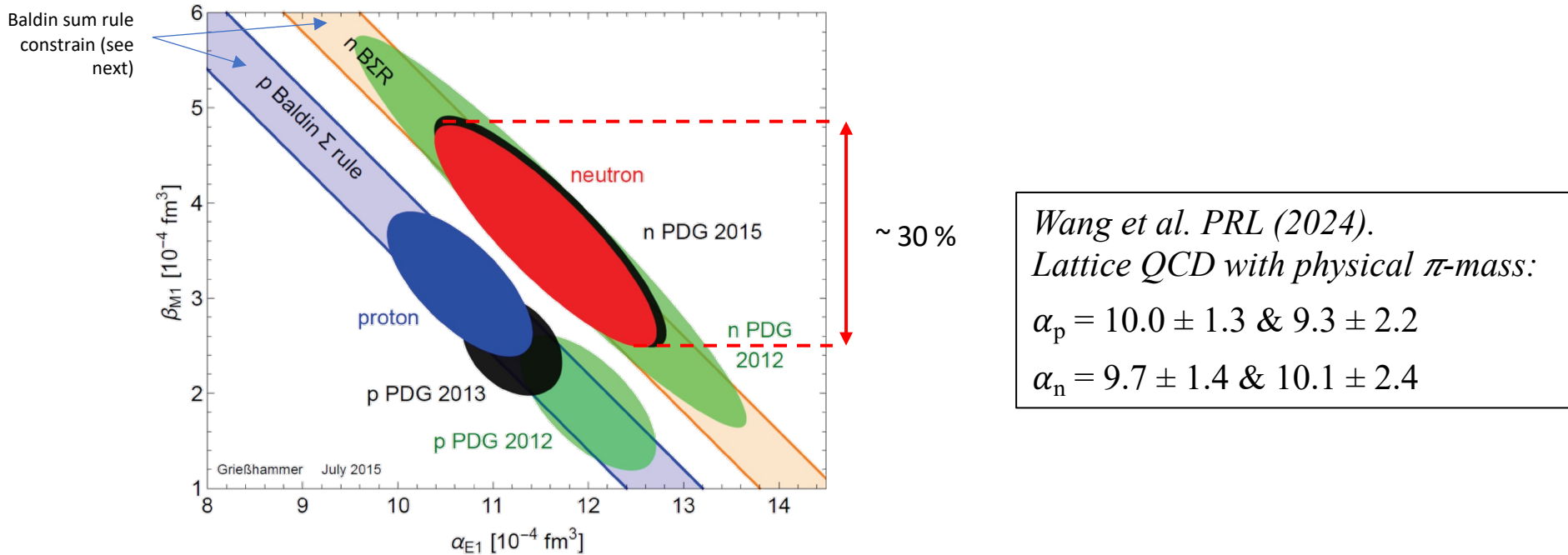
$3.15 \mp 0.35_{\text{stat}} \pm 0.2_{\Sigma} \mp 0.3_{\text{theory}}$

neutron (Baldin, NLO)
COMPTON@MAX-lab PRL 2014

$11.55 \pm 1.25_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$

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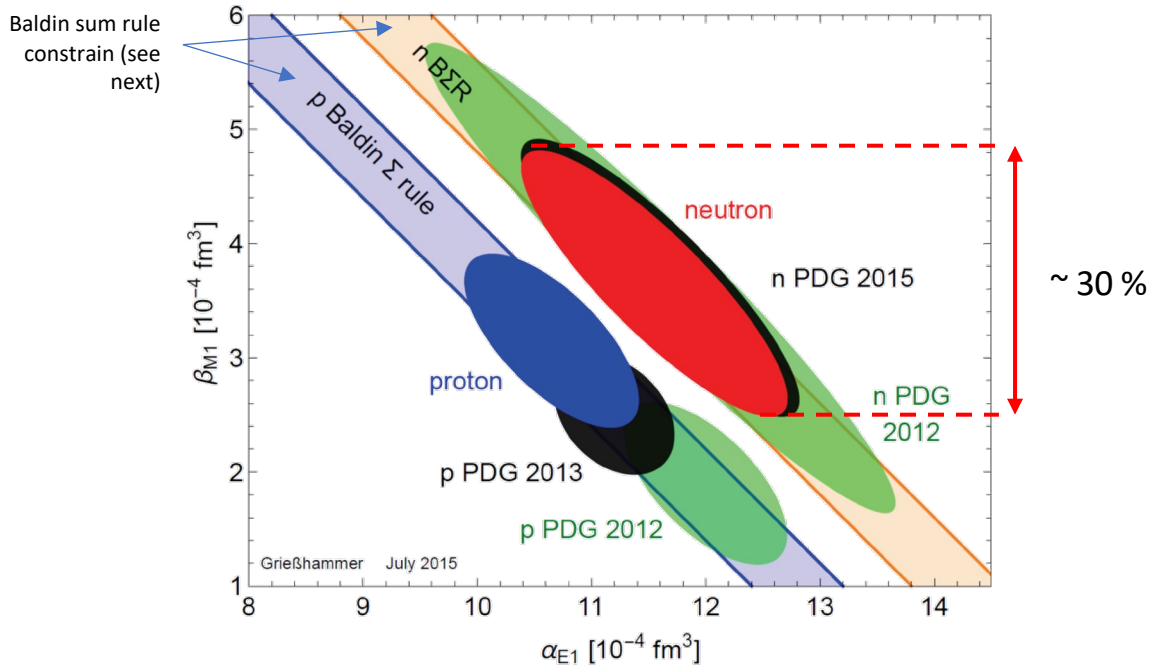
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Current knowledge of scalar polarizabilities



Wang et al. PRL (2024).
Lattice QCD with physical π -mass:
 $\alpha_p = 10.0 \pm 1.3 \text{ \& } 9.3 \pm 2.2$
 $\alpha_n = \underline{9.7 \pm 1.4} \text{ \& } 10.1 \pm 2.4$

Scalar Dipole Polarisabilities: “canonical units” $[10^{-4} \text{ fm}^3]$

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Better experimental data needed
 already to improve comparison!

proton (Baldin, $N^2\text{LO}$)
 McGovern/Phillips/hg EPJA 2013
 neutron (Baldin, NLO)
 COMPTON@MAX-lab PRL 2014

Technique: measure absolute differential cross-sections

Nuclear Compton scattering:

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{Powell} - \frac{e^2}{4\pi M_N} \left(\frac{\omega'}{\omega}\right)^2 \omega \omega' \left\{ \frac{\alpha + \beta}{2} (1 + \cos \theta)^2 + \frac{\alpha - \beta}{2} (1 - \cos \theta)^2 \right\} + O(\omega^4)$$

χ EFT theory

Cross-section of point-like particle with anomalous magnetic moment

Outgoing & incoming energy

Forward angles sensitive to $\alpha + \beta$
(cross-check with Baldin sum rule)

Backward angles sensitive to $\alpha - \beta$ (what we extract)

Baldin Sum Rule:

$$\alpha_{E1} + \beta_{M1} = \frac{1}{2\pi^2} \int_{\omega_{thr}}^{\infty} \frac{\sigma_{Tot}(\omega')}{\omega'^2} d\omega'$$

Total photo-nuclear cross-section (from other experimental data)

Technique: measure absolute differential cross-sections

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Proton polarizabilities directly from ^1H target.

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Scattering angle θ \rightarrow χ EFT theory

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Proton polarizabilities directly from ^1H target.

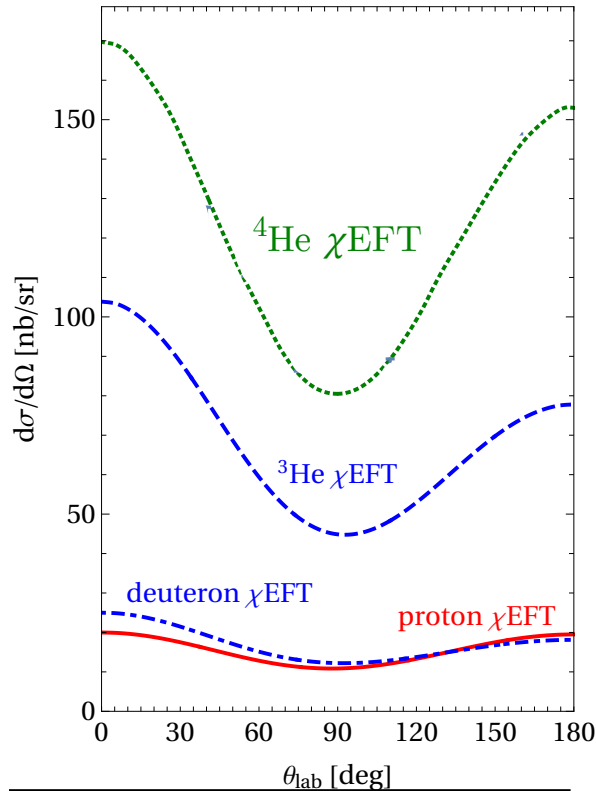
Neutron polarizabilities harder to determine:

- Uncharged, $\alpha_{E1}^n, \beta_{M1}^n$ appear at the order of ω^4 , smaller cross section
- No stable free neutron target.
- Can use light nuclear targets: **D**, **^4He** , **^3He** for different summed isoscalar combinations

χ EFT calculations of differential cross-sections

(only one energy shown) $\omega_{\text{lab}} = 60$ MeV

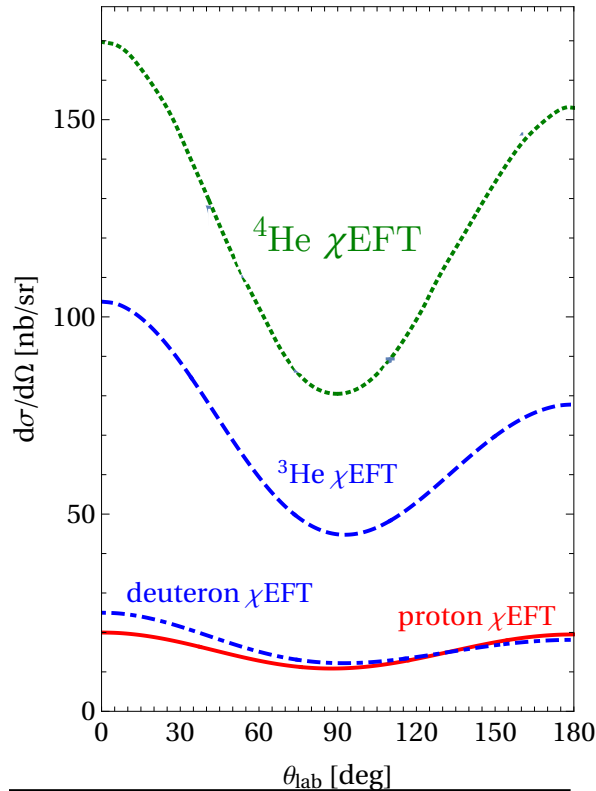
- χ EFT gives unique **angular dependence at different energies** for different α and β input values. (later, plots for different α and β for sensitivity)



(Thanks to H. Griesshammer)

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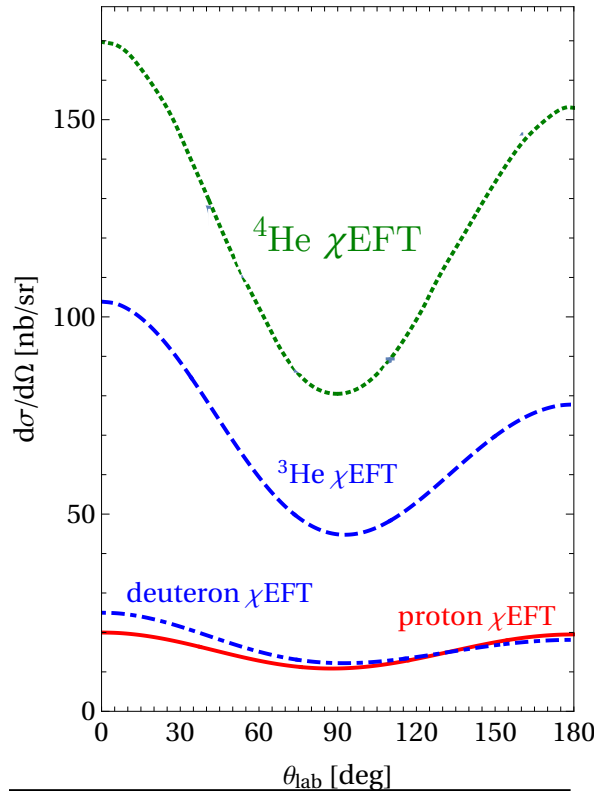


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- χ EFT gives unique **angular dependence at different energies** for different α and β input values. (later, plots for different α and β for sensitivity)
- Cross-sections are ~ 10 - 100 nb. Need **liquid target** to have high density
- Goal $\sim \pm 3$ % total uncertainty with $E_\gamma = 60 - 100$ MeV (sub pion-threshold)

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- Goal $\sim \pm 3\%$ total uncertainty with $E_\gamma = 60 - 100$ MeV (sub pion-threshold)
- While higher mass means larger cross-sections, but the theory gets more difficult and newer effects (good cross-check).
- Also need to consider systematics from **inelastic channels** (not of interest):

$$\gamma + \text{D} = \text{n} + \text{p} + \gamma' (> 2.2 \text{ MeV loss})$$

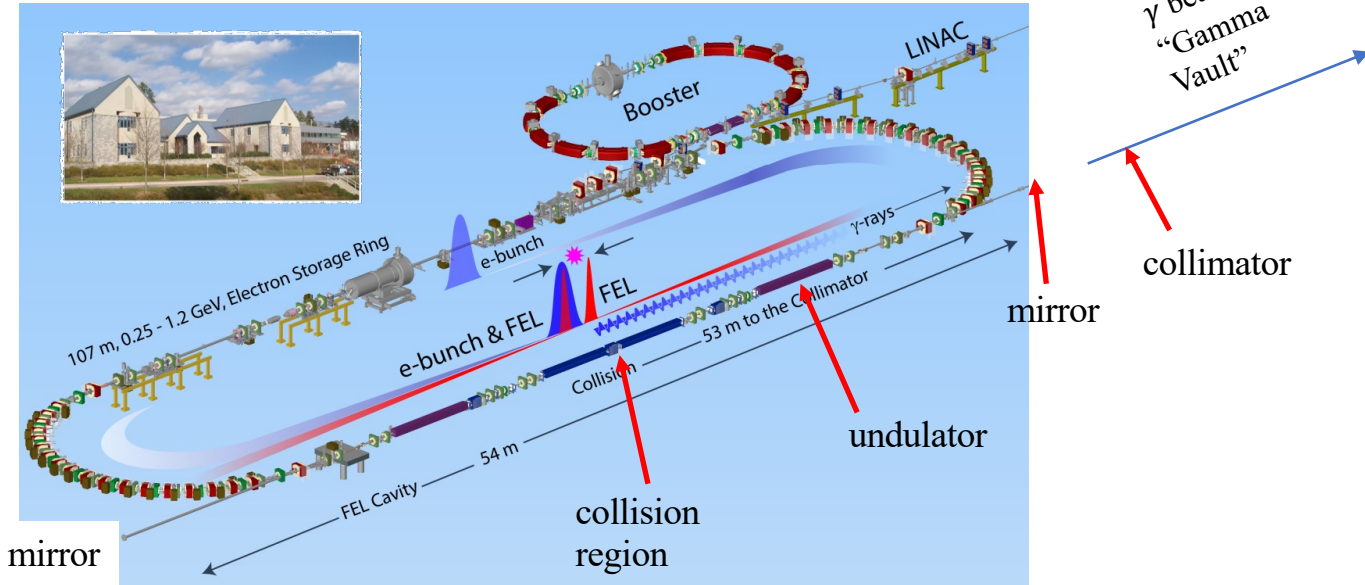
$$\gamma + {}^3\text{He} = \text{p} + {}^3\text{He} + \gamma' (> 5.5 \text{ MeV loss})$$

$$\gamma + {}^4\text{He} = \text{n} + {}^3\text{He} + \gamma' (> 20 \text{ MeV loss})$$

- Need detector resolution to separate these from (elastic) Compton scattering. Difficult for D.

High Intensity γ -Ray Source (HI γ S)

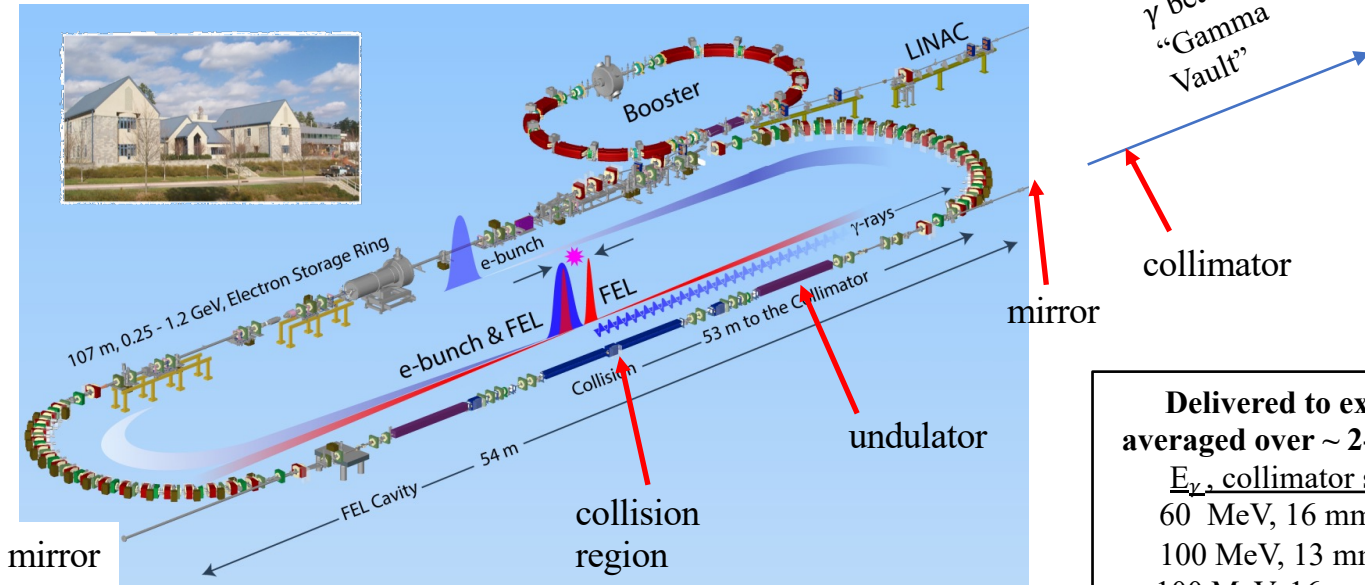
Located at Duke University & Triangle Universities Nuclear Laboratory (TUNL)



- Quasi-monoenergetic, pulsed γ -ray beams
- Synchrotron storage ring for bunches of ~ 1 GeV electrons.
- Free-electron-laser light pulse from one electron bucket reflects off mirror, and Compton scatters off next bucket of electrons to get boosted to gamma energies.
- γ beam rate and bandwidth depends on energy, collimator, mirror conditions.

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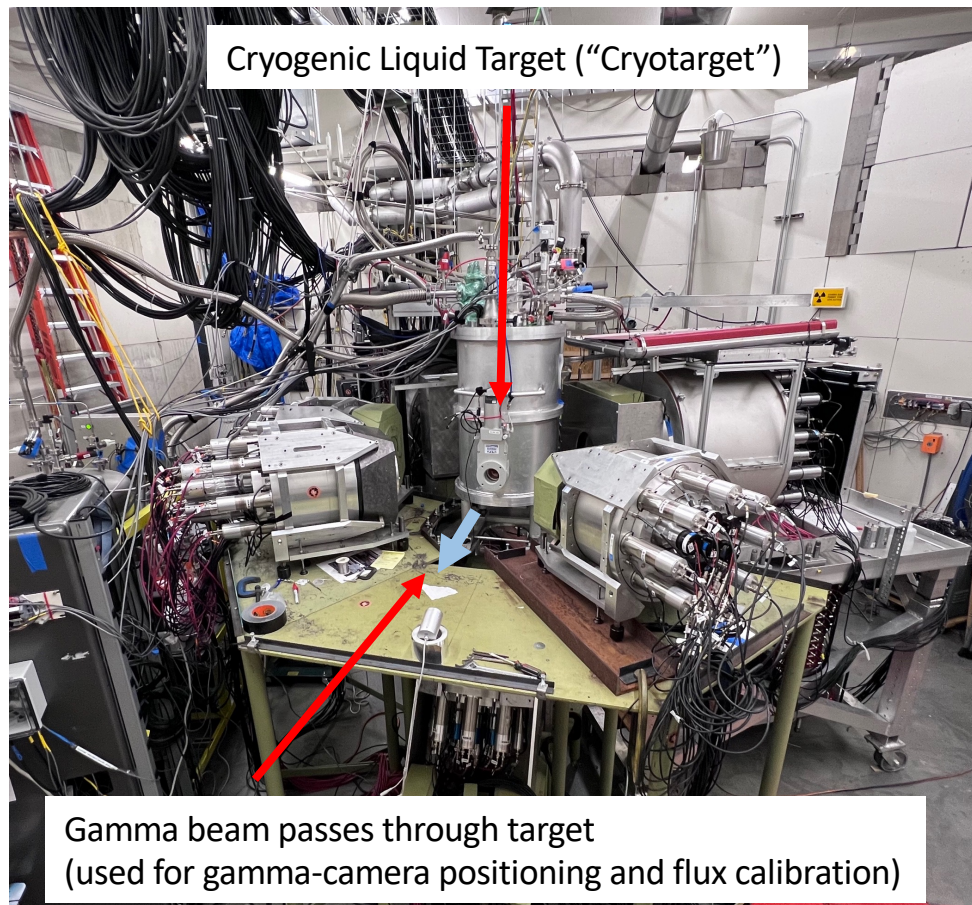
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**Delivered to experiments
averaged over ~ 2 -3 weeks each**

E_γ	collimator size	γ rate
60 MeV	16 mm	$1.4E7 \gamma/s$
100 MeV	13 mm	$4E6 \gamma/s^*$
100 MeV	16 mm	$1.2E7 \gamma/s$
87 MeV	16 mm	$1.1E7 \gamma/s$

(*1st time running @ 100 MeV)

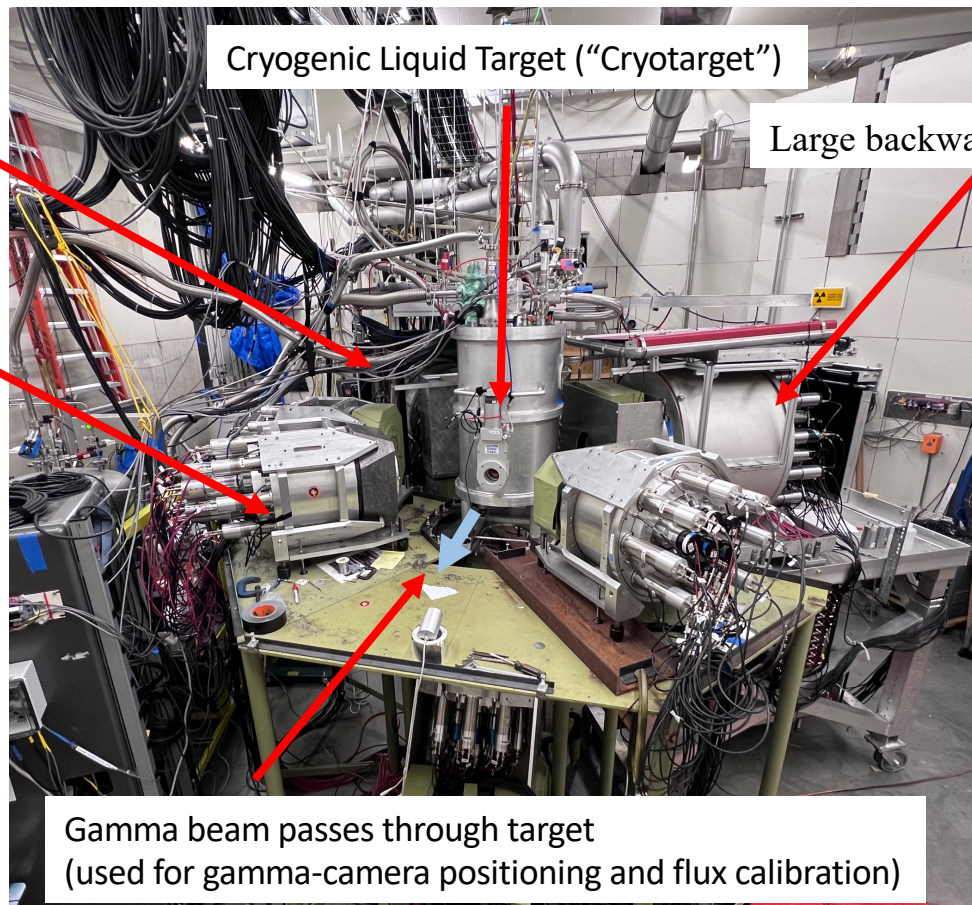
Experimental setup in Gamma Vault



Experimental setup in Gamma Vault

Large backward-angle detector #1

5x Smaller “HINDA” NaI detectors



Cryogenic Liquid Target (“Cryotarget”)

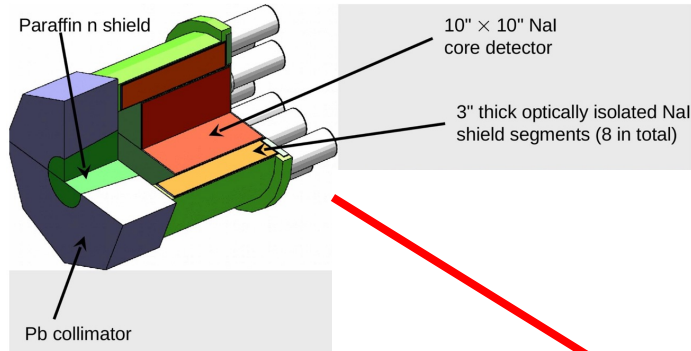
Large backward-angle detector #2

Gamma beam passes through target
(used for gamma-camera positioning and flux calibration)

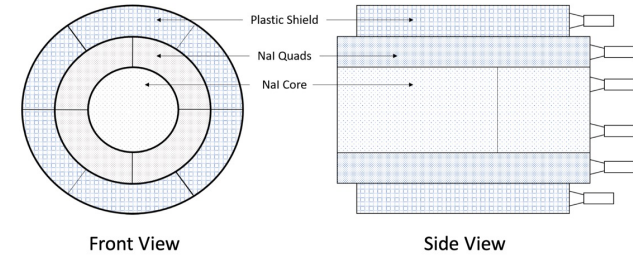
Sodium-Iodide Detectors

HINDA array (our "small" detectors)

25cm x 25 cm core



- Size needed for ~ 100 MeV γ 's
- Single crystals help with timing
- $\sim 100\%$ detection efficiency of showers
- Outer segments infer "energy leakage"

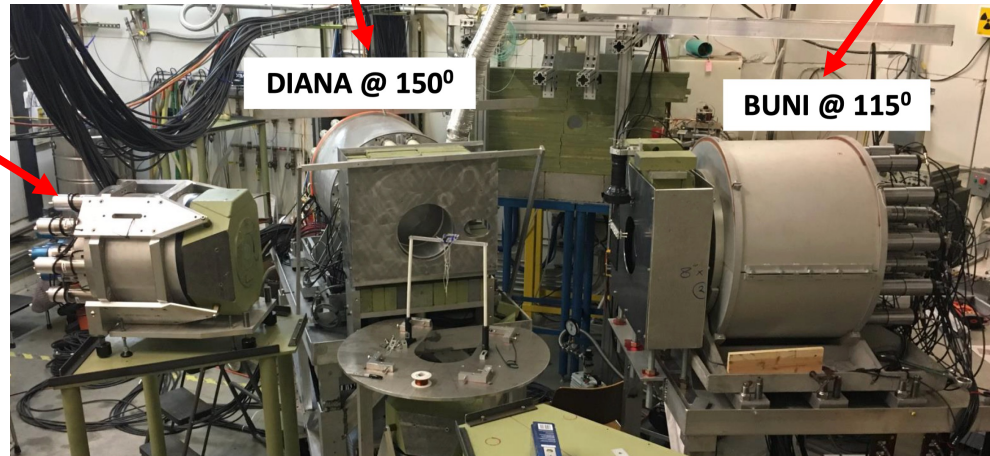


DIANA (large backward det. #1)

Core: 48-cm- \varnothing , 51-cm-long

BUNI (large backward det. #1)

Core: 2x 56-cm-long, $\varnothing 27$ cm
core glued together



Published results

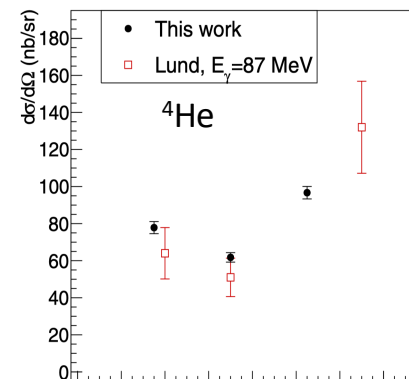
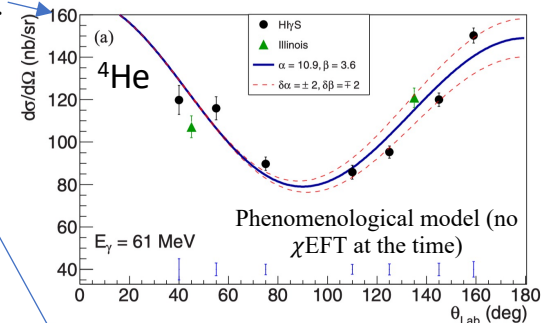
3-K cryogenic target constructed (cooling directly with 2nd stage of GM cryocooler). More on target later.

Demonstrates target, beam, and HINDA detectors.

Energy resolution to separate quasi-elastic not needed.

⁴He (60 MeV, 7 angles): Sikora et al. PRC (2017).

⁴He (80 MeV, 3 angles): Li et al. PRC (2020).



Published results

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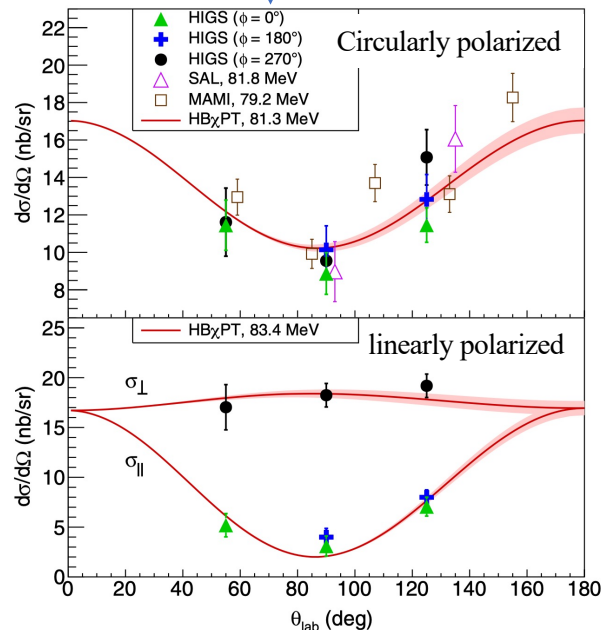
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Energy resolution to separate quasi-elastic scattering.

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⁴He (80 MeV, 3 angles): Li et al. PRC (2020).

¹H (80 MeV, 3 angles): Li et al. PRL (2022).



Our result:

$$\alpha_{E1}^p = 15.4 \pm 1.8_{\text{stat}},$$

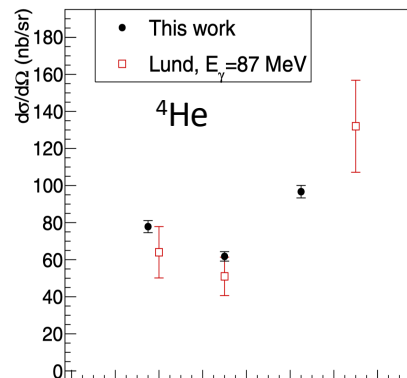
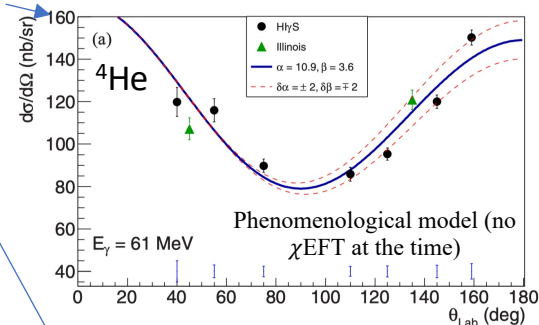
$$\beta_{M1}^p = 2.1 \pm 2.0_{\text{stat}},$$

Mornacchi et al. (MAMI), PRL (2022).
(All uncertainties added in quadrature):

$$\alpha_{E1}^p = 10.99 \pm 0.63$$

$$\beta_{M1}^p = 3.14 \pm 0.51$$

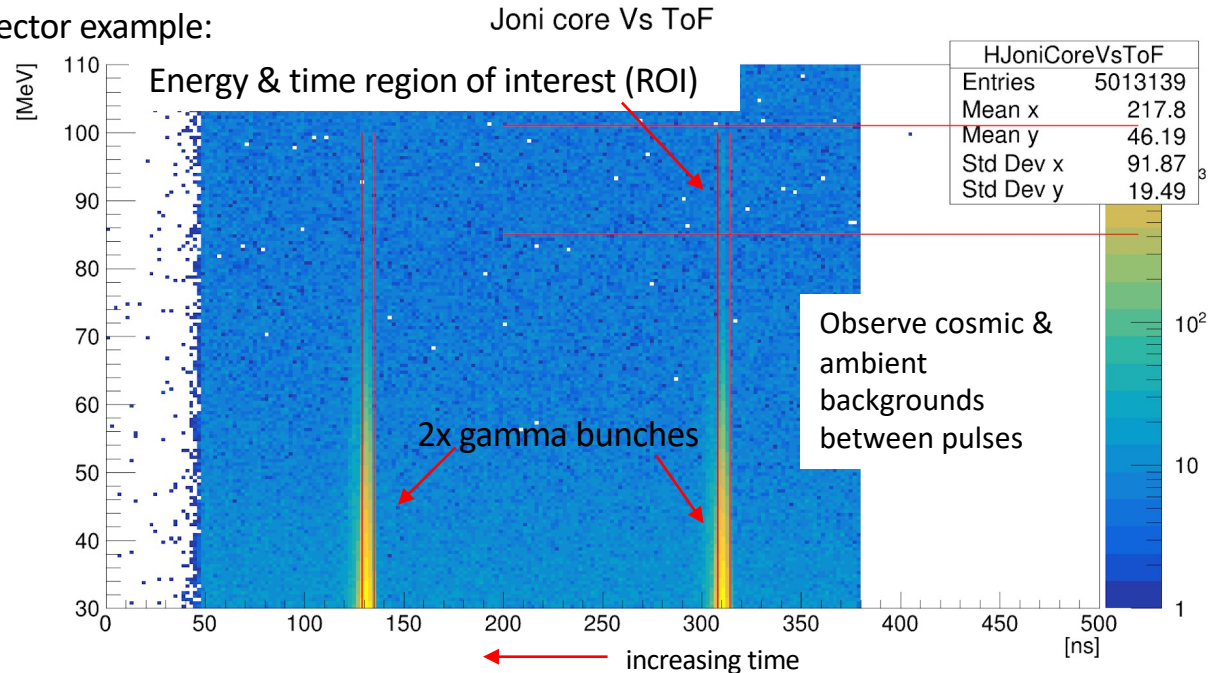
2.3 σ tension in α^p



Time and energy histograms in a detector

- Cosmic muon peak @ ~ 250 MeV, rate of tail in ROI still high \rightarrow **cosmic veto paddles**
- γ bunches (~ 10 ns wide every ~ 150 ns) allow **time cuts** and observation of beam-uncorrelated BG
- Inelastic processes lower in energy than elastic Compton \rightarrow **lower energy cut**
- Count rate in **ROI** ~ 10 - 20 /hour in large detectors
- **Empty cell subtraction** to remove beam-correlated BG due vacuum windows & cell walls + others

Smaller HINDA detector example:



Kent Leung, Nucleon Polarizabilities
Compton@HIGS, Nuclear Photonics 2025

Example data from D₂

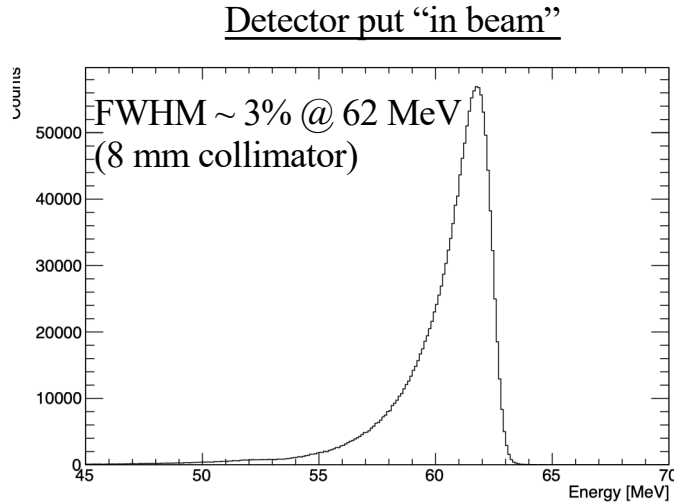


FIG. 4. Measured in-beam spectrum by DIANA. ToF and shield cuts were applied to this energy spectrum.

- Above measures beam bandwidth convoluted with detector response: $\sim [(\text{beam } 2\%)^2 + (\text{detector response } 2\%)^2]^{1/2}$

Example data from D₂

Detector put “in beam”

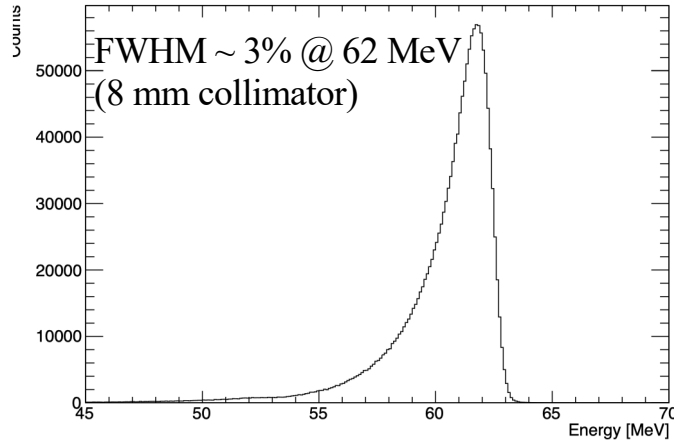
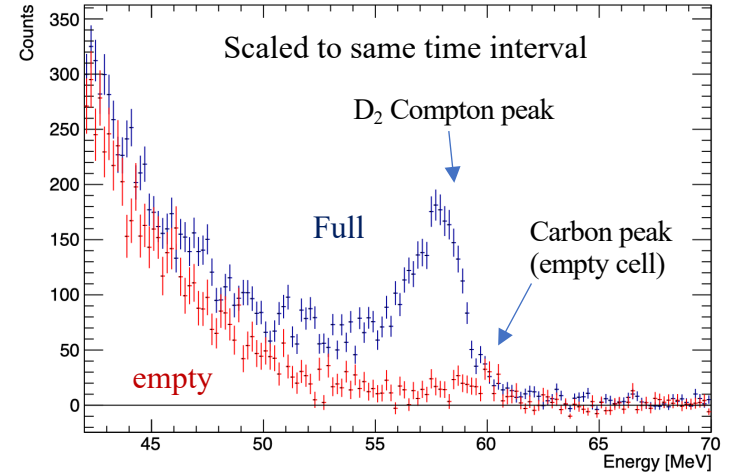


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Cell full and cell empty subtraction

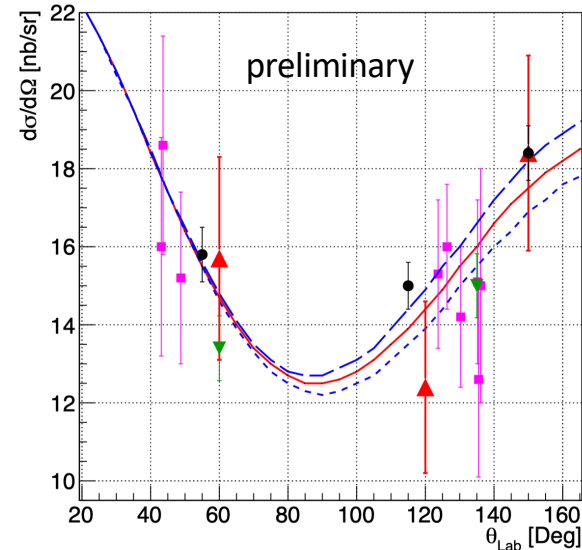
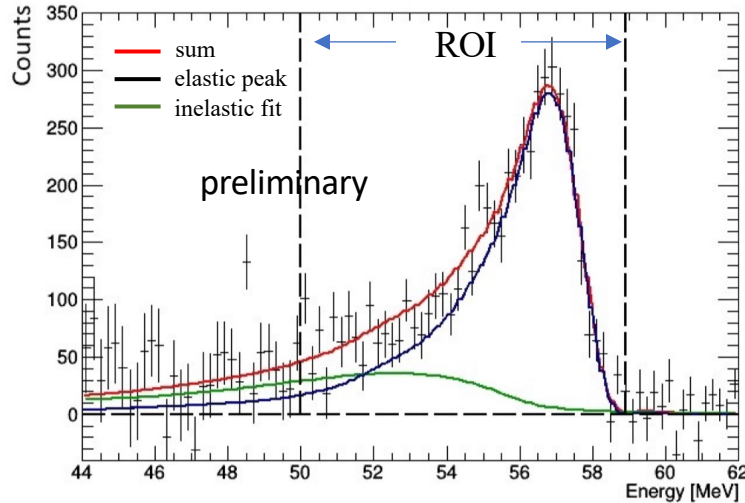


- Compton peak energy shift depends on angle of detector & mass of target.
- Corrects for BGs not from the target liquid
- Dominated by scattering by Kapton in vacuum & cell windows (carbon with high Z)

Deuteron Compton Scattering

- 60 MeV on deuteron data taken in 2022 (**Danula Godagama, PhD thesis**). Paper in draft.
- Needed the large DIANA & BUNI detectors (at backwards angles) to separate inelastic contribution
- **Adding only our 3x $d\sigma/d\Omega$ to the world data set of ~ 50 values to a χ EFT fit, we improve α_n & β_n by $\sim 10\text{-}15\%$**

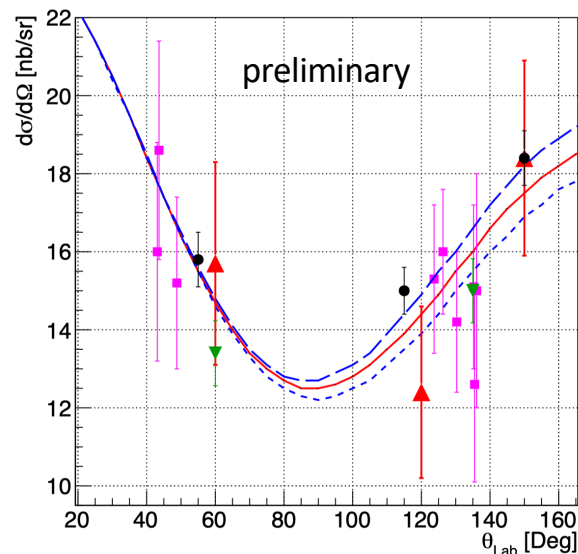
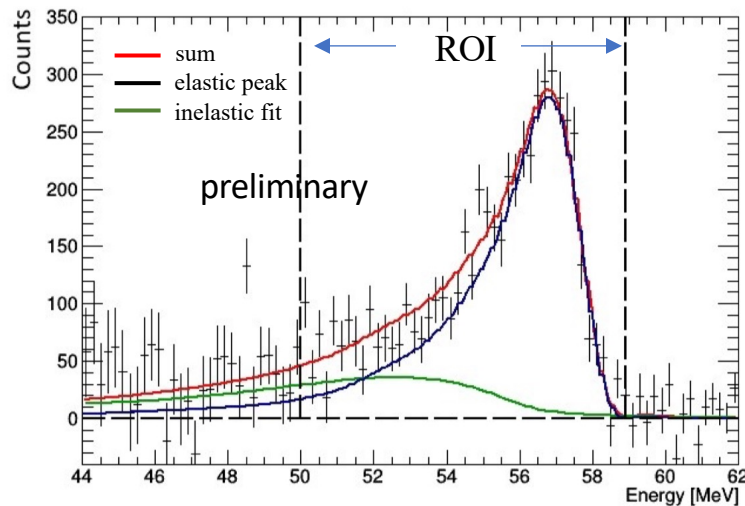
Example: 150° HINDA detector.



Deuteron Compton Scattering

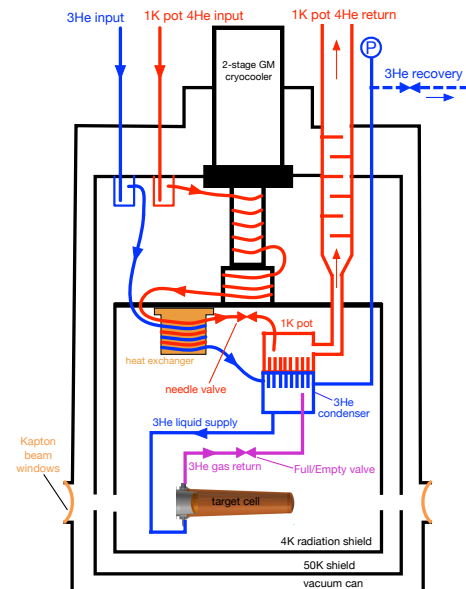
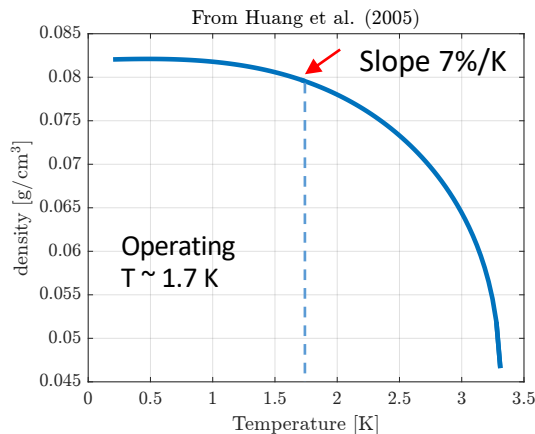
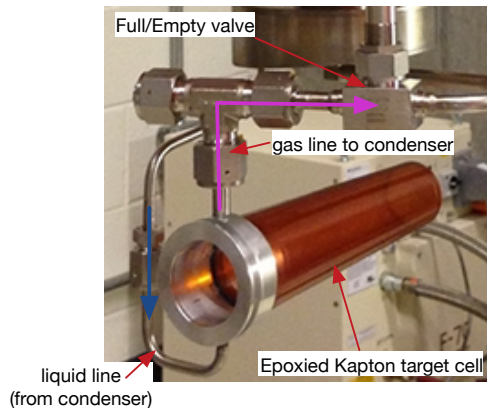
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- Besides statistics, still contending with sizable correction for inelastic falling into ROI.
- ${}^3\text{He}$ (never done before!) balances good statistics, good theory (mass 3), and **inelastic 5 MeV away**.

Example: 150° HINDA detector.



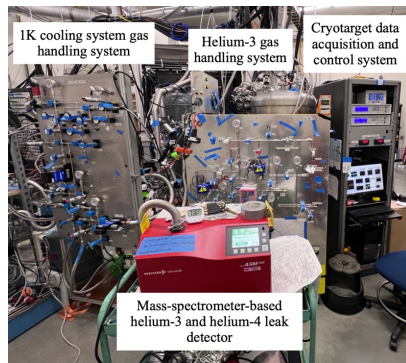
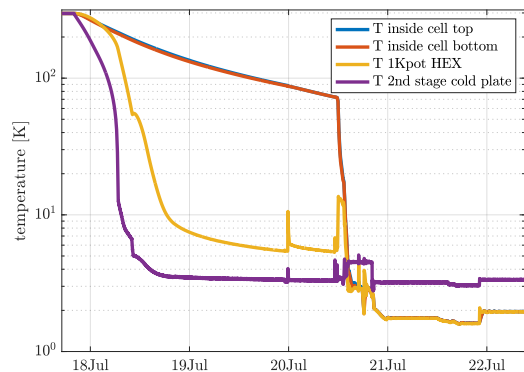
L³He Cryotarget

- Kapton cell & vacuum window to reduce backgrounds (and for optical access). Made from 0.1-mm-thick epoxied-together sheets. (Flimsy!)
- 1.7 K operation temperature for density & density stability
- In-house recirculating ⁴He 1K pot precooled with 1.5 W @ 4K GM cryocooler.
- Target liquid cooled with a thermosiphon loop
- 0.3 L of liquid in cell -> 300 bar-L of ³He. (~50 bar-L in vapor above liquid.)



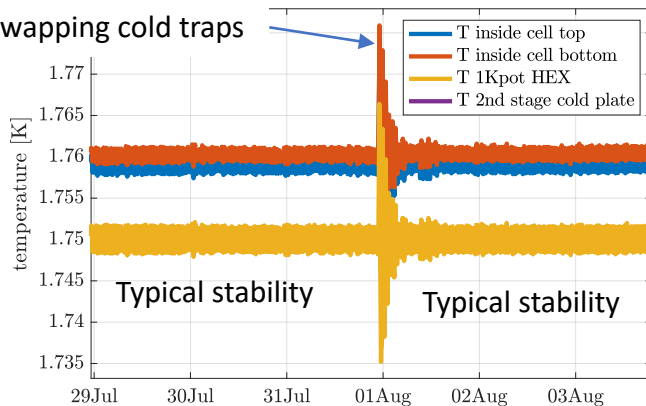
Cryogenic performance

Initial cooldown

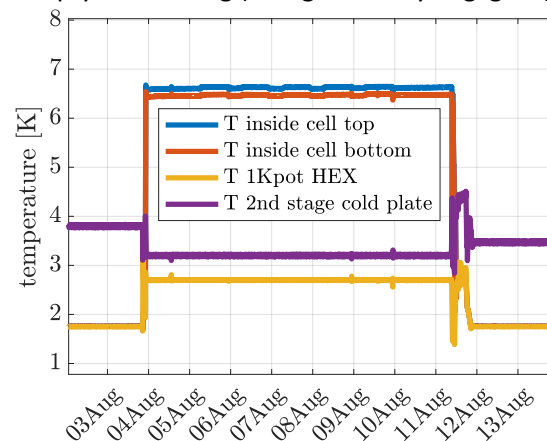


- Ran 1st time Summer '24.
- 2x 4 weeks continuously cold runs
- Temperature stability at operating better than ± 3 mK
- Did not observe any loss of ^3He
- $< 1 \text{ bar}\cdot\text{L}$ used for RGA studies
- ^4He contamination in $^3\text{He} \sim 2\%$ level

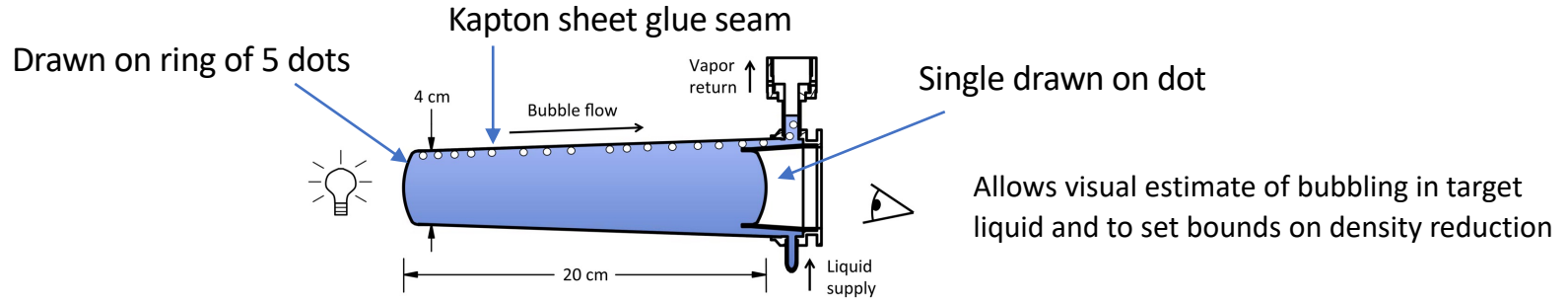
Planned procedure swapping cold traps



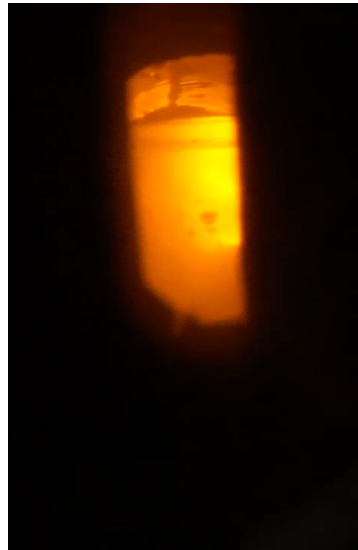
Empty cell running (cold gas density negligible)



Visual access to target cell



During condensing and filling cell with ^3He



During operation when full
(Light illuminating fluid,
which heats it, turned off
during production data)



Systematics from ^3He cryotarget

- Summer 2024: 1st experiment @ 60 MeV to test system (**Ethan Mancil PhD thesis**)
- Main physics run: 100 MeV, 360 hrs production data (~12hr /day), $\pm 3\text{-}4\%$ statistics reached (**Jingyi Zhou PhD thesis**)

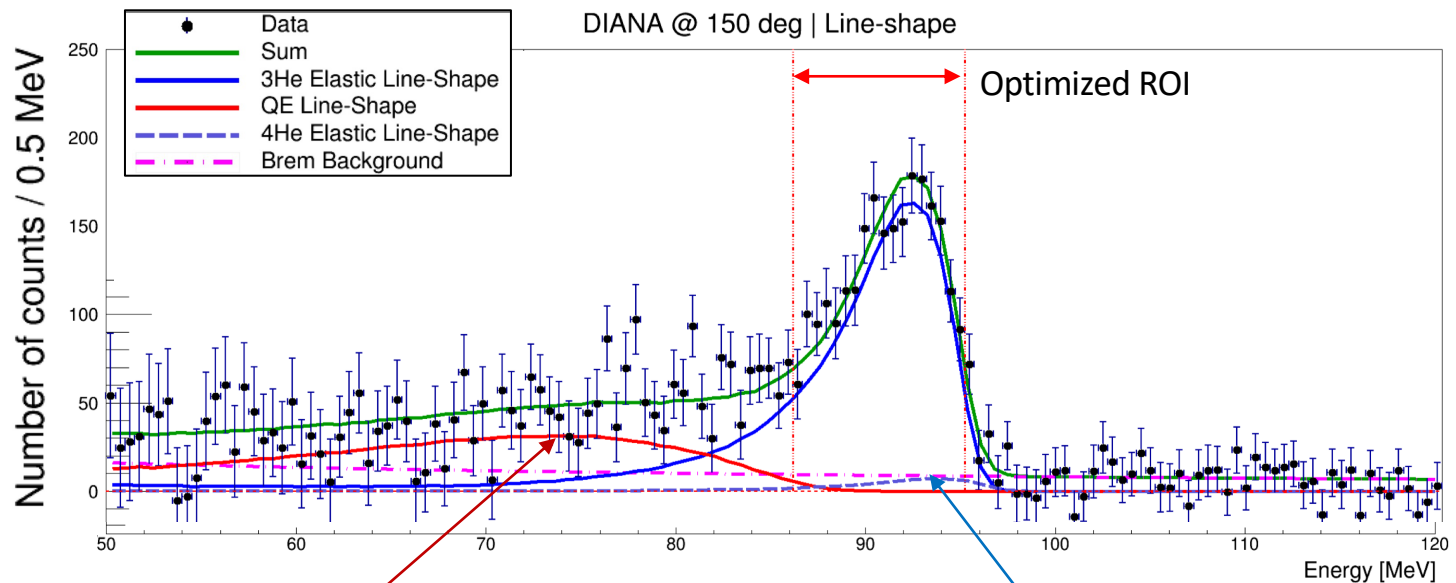
From ^3He cryotarget
design document:

Source of uncertainty	Uncertainty
Target length measurement at room temperature	0.2 %
Thermal contraction	0.1 %
Pressurized cell flexing against vacuum	0.8 %
Thermometer uncertainty	< 0.1 %
Temperature stability	< 0.1 %
Bubbling	< 0.1 %
Temperature gradients	0.7%
Total (added in quadrature)	1.1 %

Measured ^3He
cryotarget
performance from
Summer 2024:

- **Temperature stability:** better than ± 5 mK leads to $< 0.1\%$ systematic
- **Temperature gradient:** measured < 10 mK difference top and bottom. Systematic $< 0.1\%$
- **Thermometer calibration uncertainty:** $< \pm 10$ mK. Systematic $< 0.1\%$ (radiation levels are low)
- **Bubbling:** analyze videos and develop theory (depends on latent heat, bubble velocity, etc.)
- **Length of target** (pressure flexing and beam position on curved end windows): under study
- **^4He contamination:** Measured $\sim 2\%$ ^4He in ^3He gas inventory. ^4He is soluble in L^3He . RGA studies.
- **Note:** literature values for mass density of liquid ^3He is $\pm 1\text{-}2 \%$

^3He @ 100 MeV results: DIANA 150° detector



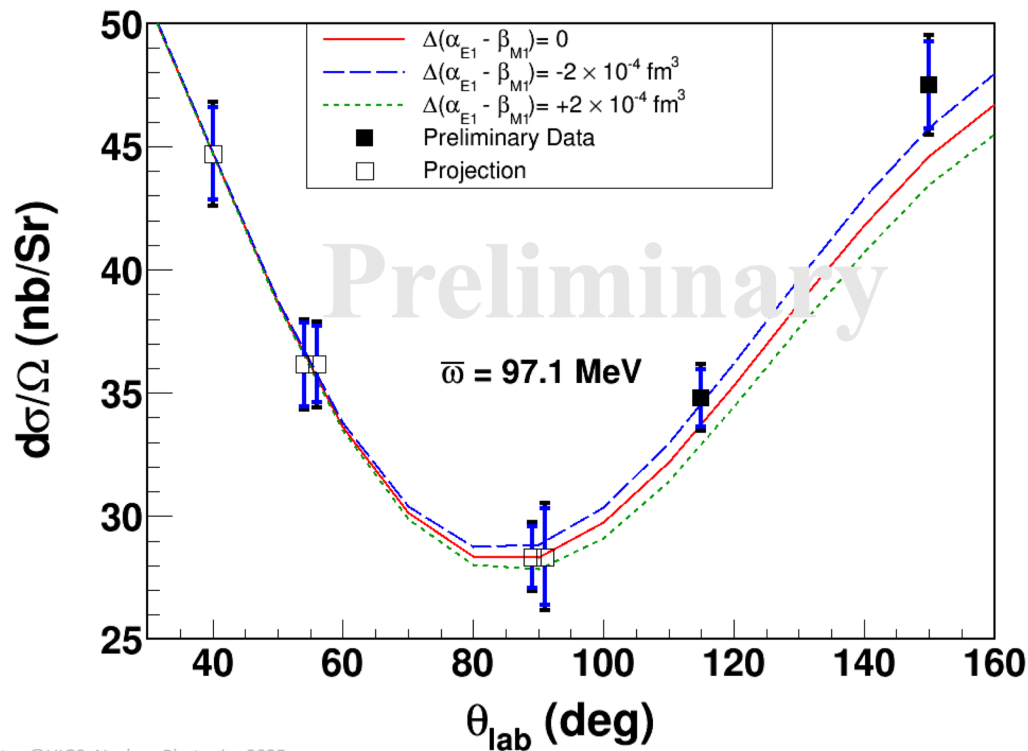
Better separation of inelastic scattering c.f. D_2

^4He Compton contamination. ~2% of ^4He determined by RGA & cross-section from measurements & theory (precision not needed).

(Thanks, Jingyi Zhou for this and next figure.)

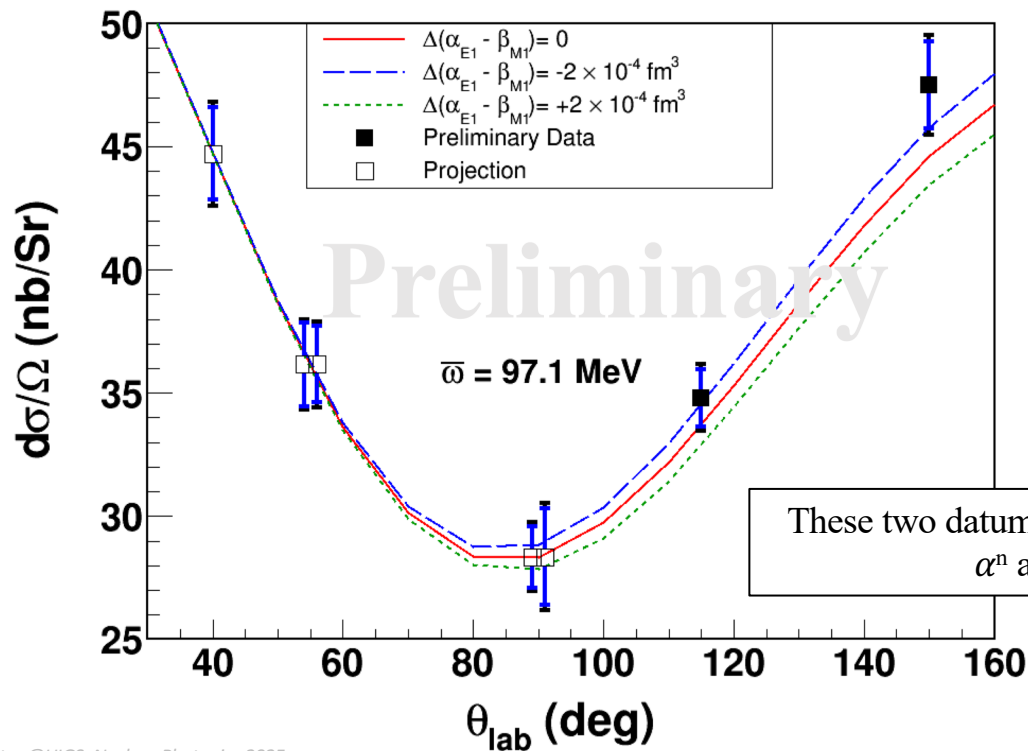
^3He results: differential cross-sections

- Recall: No other data to compare because first time this process has been measured!
- Large backward angles constrain physics.
- Forward angles verifies the Baldin sum rule.
- ("Projection" = still working on analysis of data from these detectors.)



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Compton@HIGS summary & future

- **^1H @ 80 MeV:** 2.3 σ tension in proton polarizabilities. Proposing to revisit at 100 MeV
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- D @ 60 MeV (3 datum) improved **neutron polarizabilities by 10-15%**. Proposing to revisit at 80 MeV.
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- This summer (2025) ran with **^4He @ 90 & 100 MeV** (Mitchell Lewis PhD thesis). There is now χEFT theory for ^4He . Early estimates $\sim 30\%$ improvement in neutron polarizabilities.
- R&D for next generation of experiments: **spin polarizabilities @ HIGS** below pion threshold.

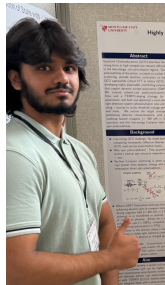
Spin polarizabilities @ HIGS

- Compton amplitude $\mathcal{O}(\omega^3)$ gives 4x **spin polarizabilities**: characterizes stiffness of nucleons' spin degrees of freedom to photons
- Requires **polarized beams + polarized target** \rightarrow **Dynamic Nuclear Polarized (DNP)** *protons* in polymers
- To suppress backgrounds from carbon scattering \rightarrow scintillating target

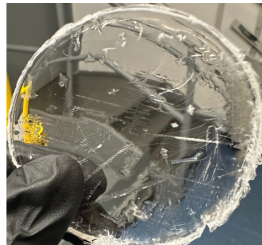
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- DNP: **polarize electrons in ~ 5 T**. Apply ~10 mW microwaves @ ~0.7 K to **transfer e-polarization to nuclear pol.**
- **“Frozen spin”**: remove 5 T magnet (blocks outgoing photons), **cool down to 10 mK**, and leave in small ~ mT coils
- Ben van den Brandt (PSI) was developing **TEMPO (free-radical) doped scintillating films**
- **R&D launched** at Montclair

Undergrad
Taha Qadir



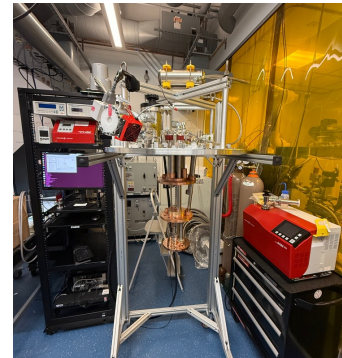
solvent-cast polyvinyl
toluene scintillating films



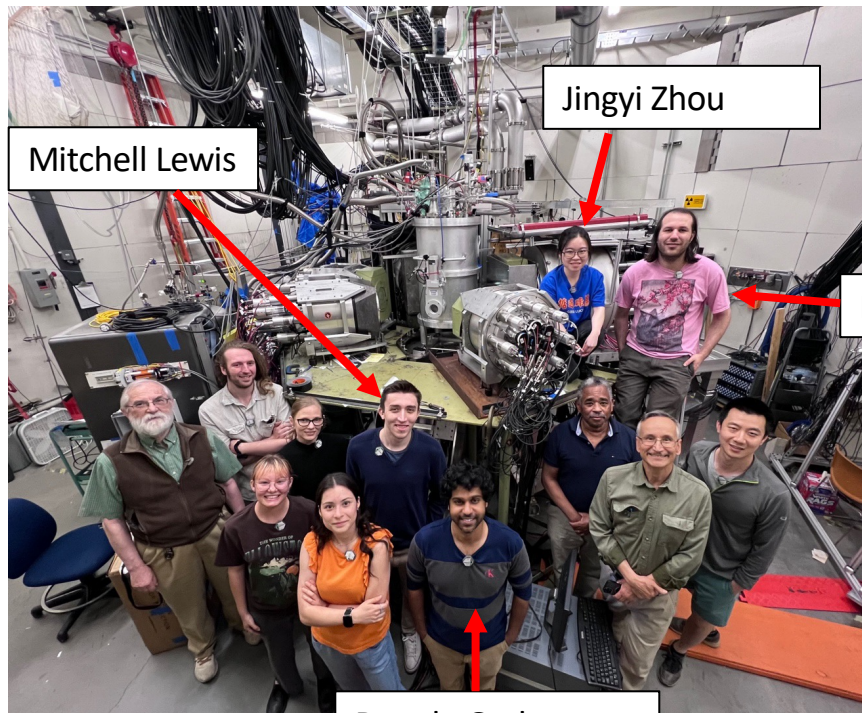
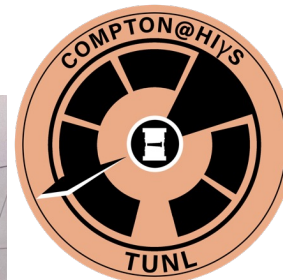
scintillation tests
with PMT



0.4 K dry-³He fridge @ Montclair



Thank you!



PhD thesis projects

Danula Godagama