

Electromagnetic response functions and electric dipole polarizability from coupled-cluster theory



Francesco Marino



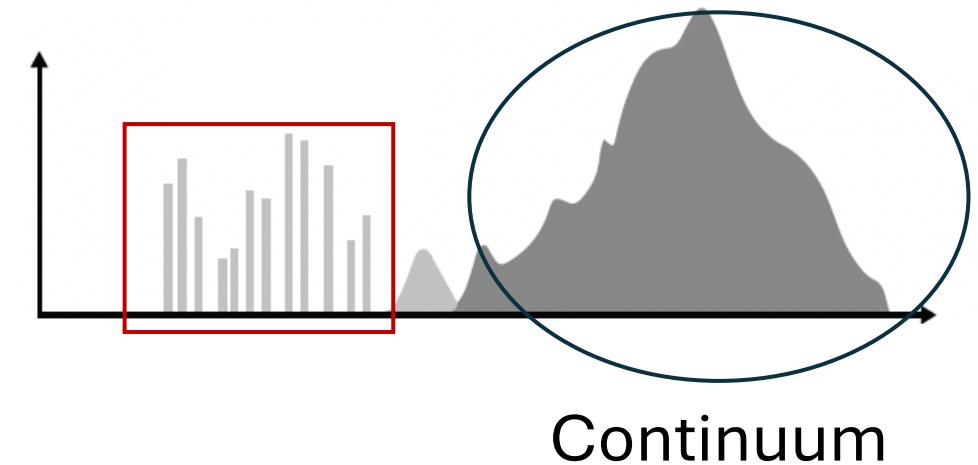
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Johannes Gutenberg-Universität Mainz



Nuclear response functions

$$R(\omega) = \sum_f |\langle \Psi_f | \Theta | \Psi_0 \rangle|^2 \delta(E_f - E_0 - \omega)$$

Θ : excitation operator



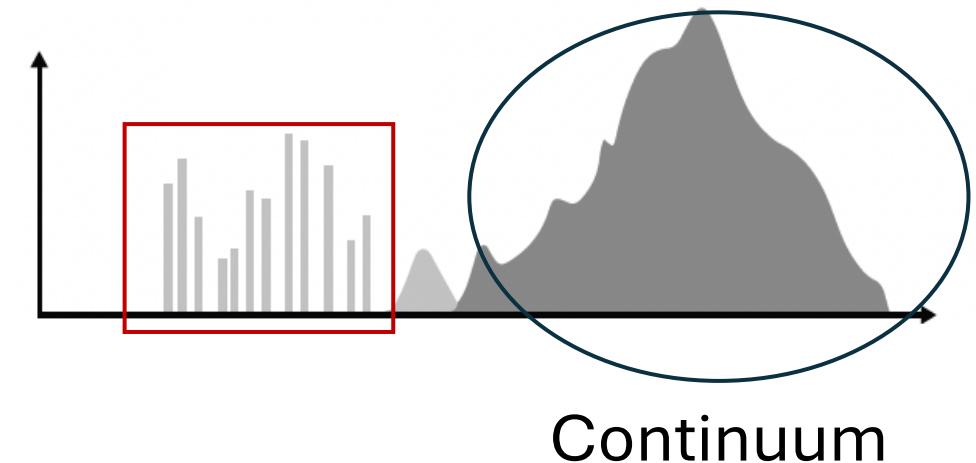
Nuclear response functions

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Θ : excitation operator

Electric dipole field

$$\Theta = \sum_{i=1}^z (\mathbf{r} - \mathbf{r}_{cm})_3$$



The response to the electric dipole is related to the photo-absorption cross section

$$\sigma_\gamma(\omega) = 4\pi^2 \alpha \omega R(\omega)$$

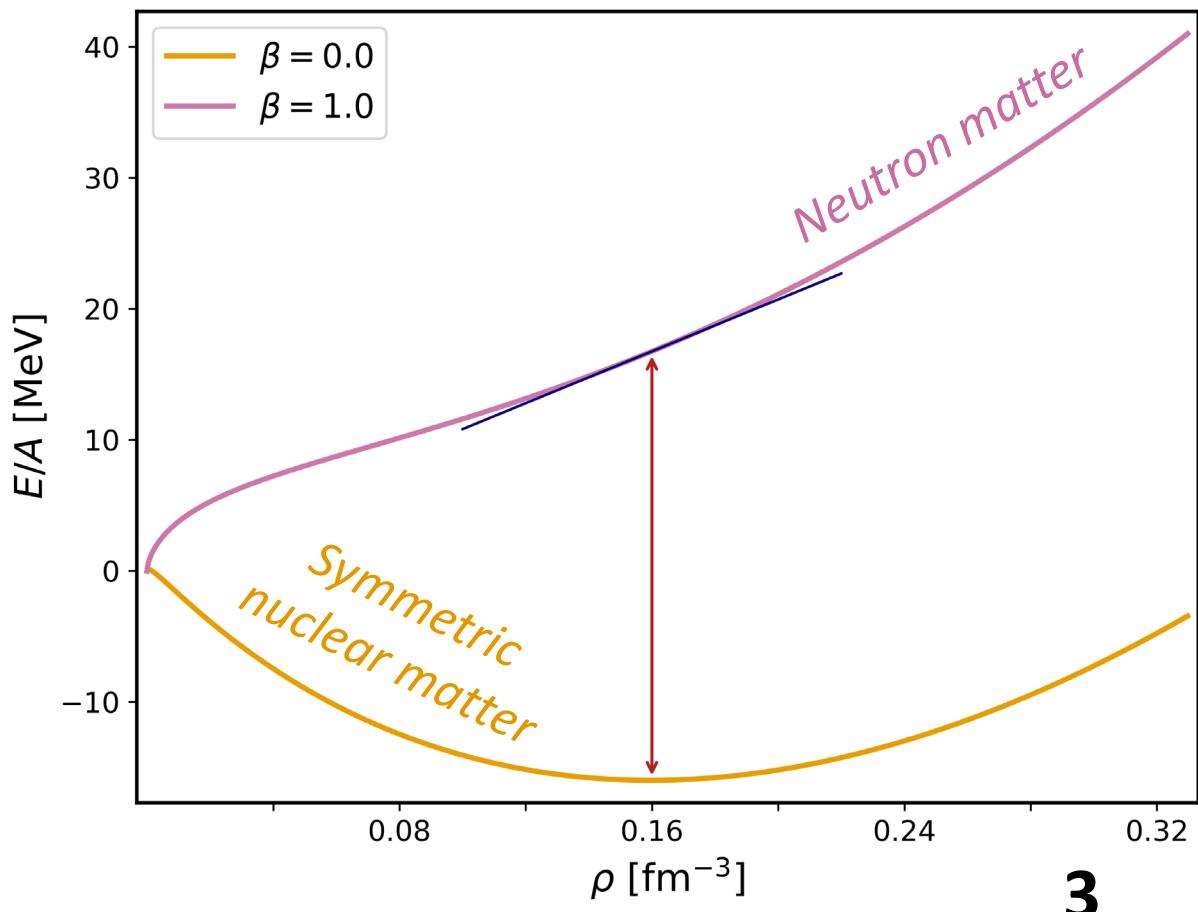
The equation of state of infinite nuclear matter

$$\frac{E}{A}(\rho, \beta) = \frac{E}{A}(\rho, 0) + S(\rho)\beta^2 + O(\beta^4)$$

$$\rho = \rho_n + \rho_p$$
$$\beta = \frac{\rho_n - \rho_p}{\rho}$$

Symmetry energy

$$S(\rho) = J + \frac{\rho - \rho_0}{3\rho_0} L$$



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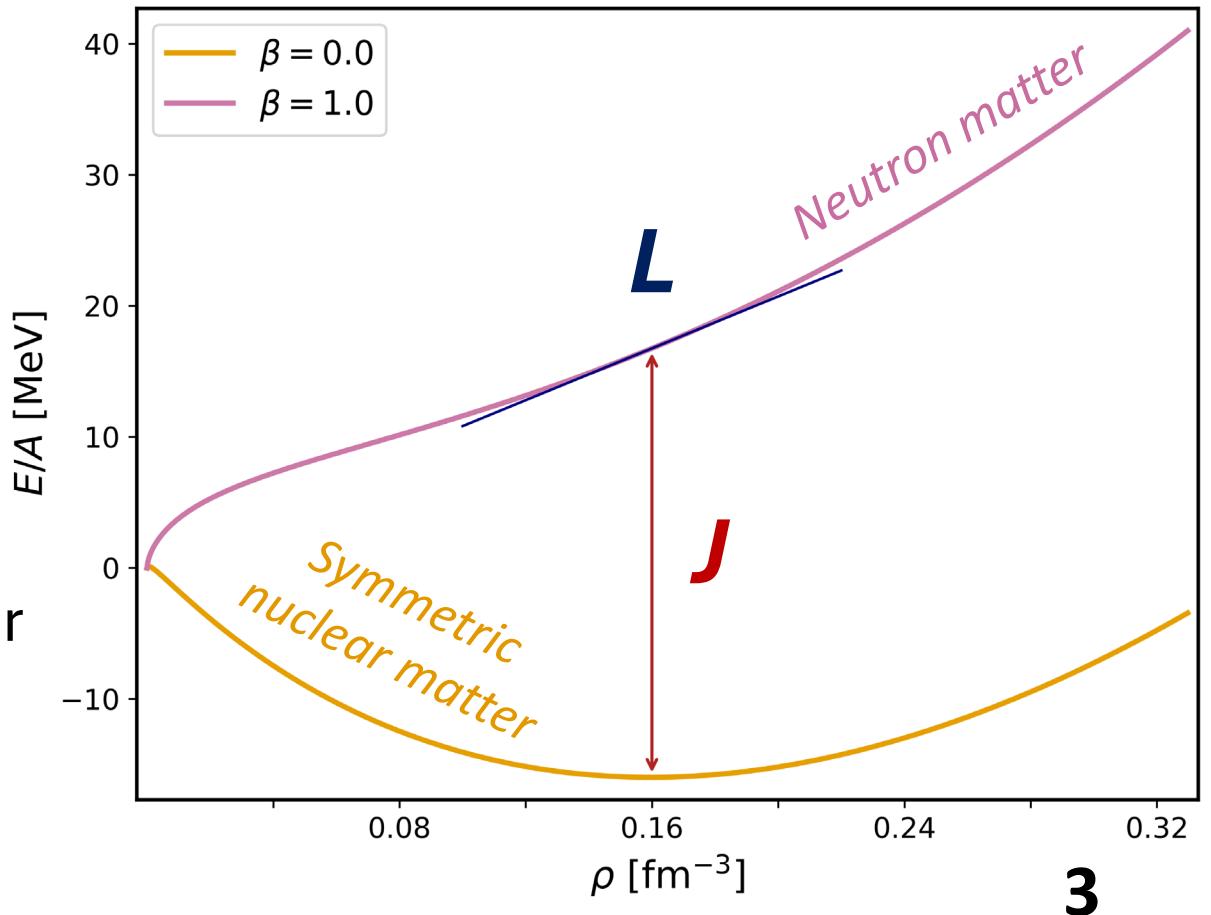
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Symmetry energy

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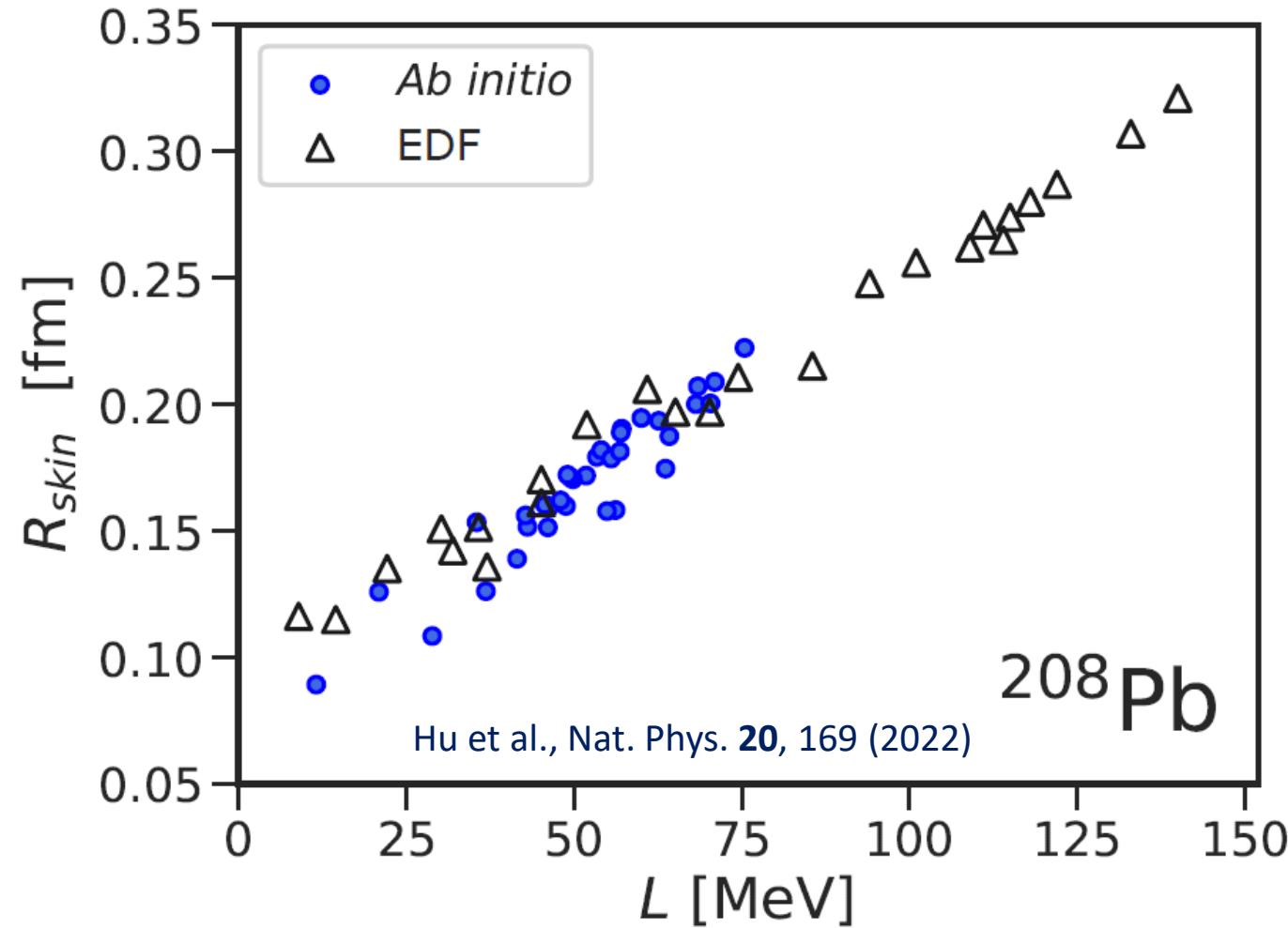
J: symmetry energy
at saturation density

L: slope parameter



How do we constrain the slope L ?

Exploit correlations with measurable nuclear properties



Neutron skins

$$R_{skin} = R_n - R_p$$

Well-defined trend, but difficult to measure neutron radii

The electric dipole polarizability

Electric dipole polarizability $\alpha_D = 2\alpha \int d\omega \frac{R(\omega)}{\omega}$

Neumann-Cosel et al.,
Front. Phys. **13**, 1629987 (2025)

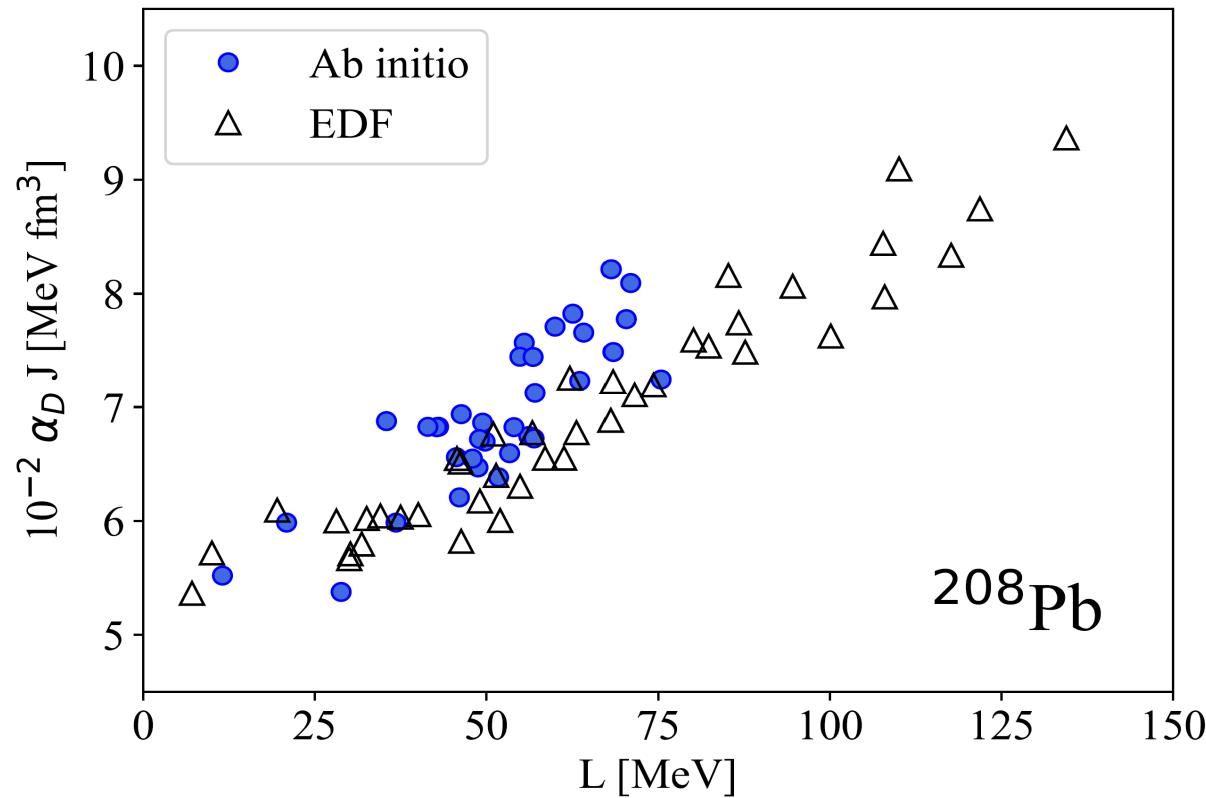
The electric dipole polarizability

Electric dipole polarizability

$$\alpha_D = 2\alpha \int d\omega \frac{R(\omega)}{\omega}$$

Neumann-Cosel et al.,
Front. Phys. **13**, 1629987 (2025)

Courtesy: Francesca Bonaiti



Strong linear correlation
between α_D and the slope
of the symmetry energy L

Coupled-cluster theory

Coupled-cluster (CC) ground-state

$$|\Psi_0\rangle = e^T |\Phi_0\rangle$$

Coupled-cluster theory

Coupled-cluster (CC) ground-state

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$|\Phi_0\rangle$ Reference state

Coupled-cluster theory

Coupled-cluster (CC) ground-state

$$|\Psi_0\rangle = e^T |\Phi_0\rangle$$

$|\Phi_0\rangle$ Reference state

CCSDT-1: approximate triples

$$T \approx \boxed{T_1 + T_2 + T_3} + \dots \quad n\text{-particle } n\text{-hole amplitudes}$$

CCSD

Coupled-cluster theory

Coupled-cluster (CC) ground-state

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CCSDT-1: approximate triples

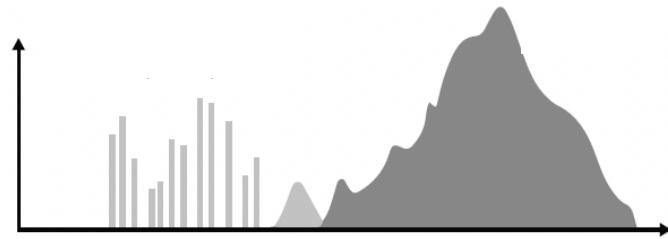
$$T \approx T_1 + T_2 + T_3 + \dots \quad n\text{-particle } n\text{-hole amplitudes}$$

CCSD

$$|\Psi_0\rangle = \begin{array}{c} \text{---} \\ | \end{array} \dots + \dots$$

$|\Phi_0\rangle$ $1p1h$ excitations $2p2h$ excitations

The LIT-CC method



$$R(\omega) = \sum_f |\langle \Psi_f | \Theta | \Psi_0 \rangle|^2 \delta(E_f - E_0 - \omega)$$

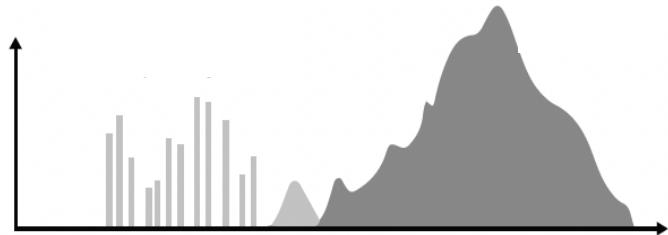
Lorentz integral transform (LIT)

$$L(\sigma, \Gamma) = \frac{\Gamma}{\pi} \int d\omega \frac{R(\omega)}{(\omega - \sigma)^2 + \Gamma^2} = \frac{\Gamma}{\pi} \langle \tilde{\Psi}_L | \tilde{\Psi}_R \rangle$$

Bound pseudo-state $|\tilde{\Psi}_R\rangle$

Bacca et al, Phys. Rev. Lett. **111**, 122502 (2013)

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✓ Sum rules of the response directly from $L(\sigma, \Gamma \rightarrow 0)$

≡ Inversion needed to find $R(\omega)$

The LIT-CC method

Pseudo-states from Schödinger eqs with source term

$$(\bar{H} - z)|\tilde{\Psi}_R\rangle = \bar{\Theta}|\Phi_0\rangle \quad z = \sigma + i\Gamma$$

Similarity-transformed operators

$$\begin{aligned}\bar{H} &= e^{-T} \hat{H}_N e^T \\ \bar{\Theta} &= e^{-T} \hat{\Theta}_N e^T\end{aligned}$$

Correlations included through ground-state amplitudes T

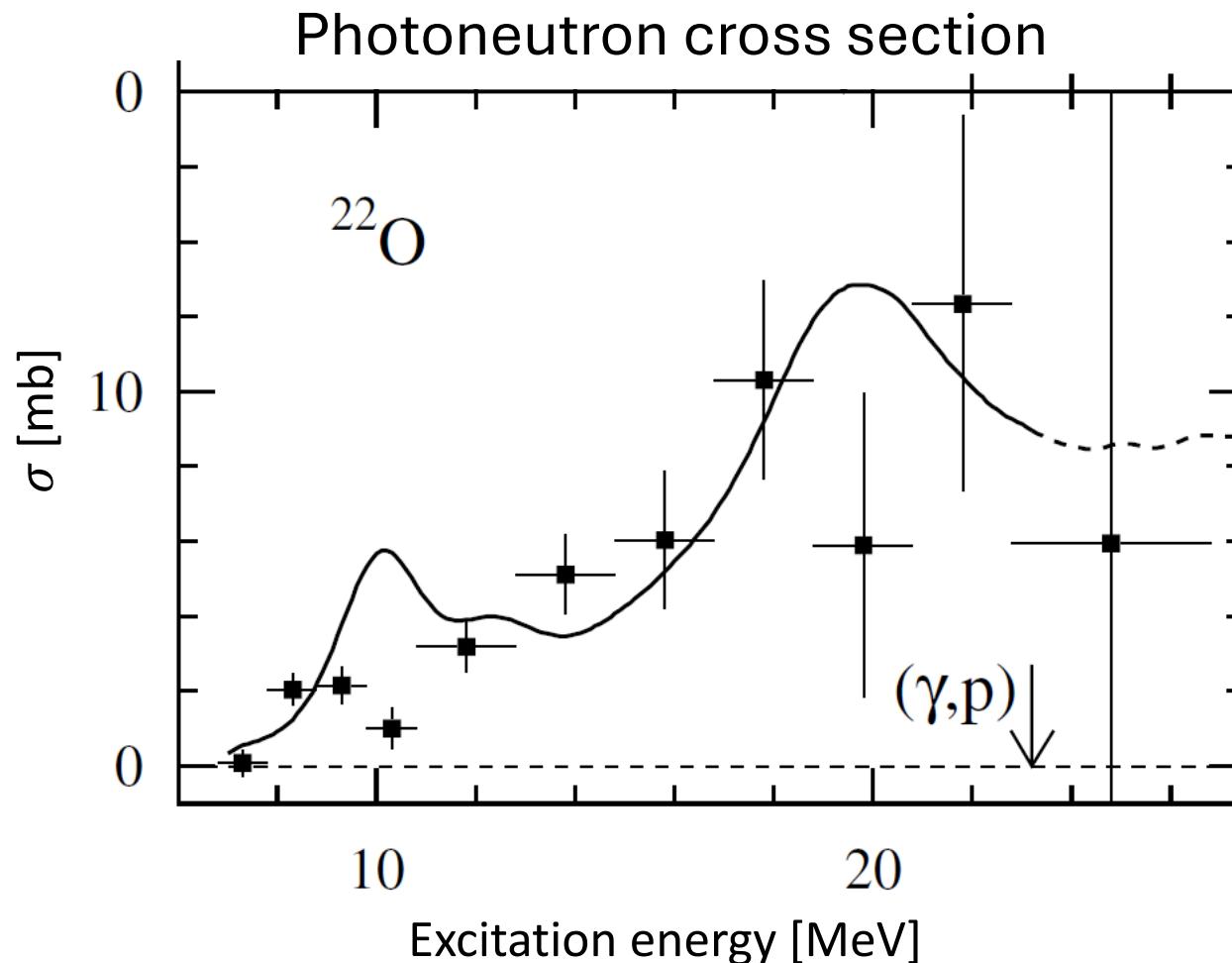
$|\tilde{\Psi}_R\rangle$ obtained in the space of 2p2h configurations

Electric dipole response of ^{22}O from coupled-cluster theory

A study of ^{22}O

^{22}O is one of the few neutron-rich nuclei for which the dipole response has been measured

Leistenschneider et al., Phys. Rev. Lett. 86, 5442 (2001)

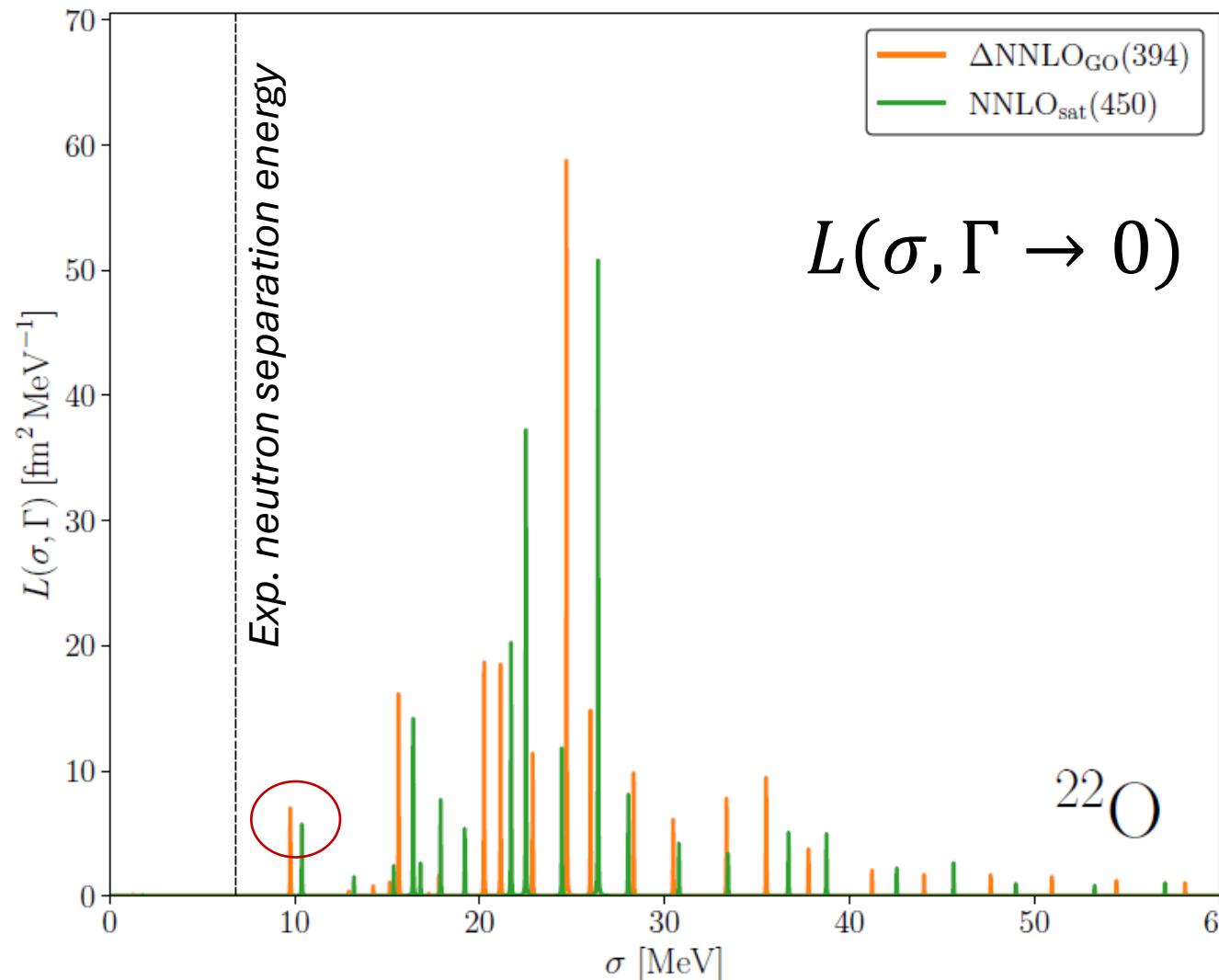


Photoneutron cross-section extracted
from electromagnetic excitations
cross-sections

Caveats: limited energy range, proton
decay channel not resolved

Theory predictions: the LIT

Discretized response functions

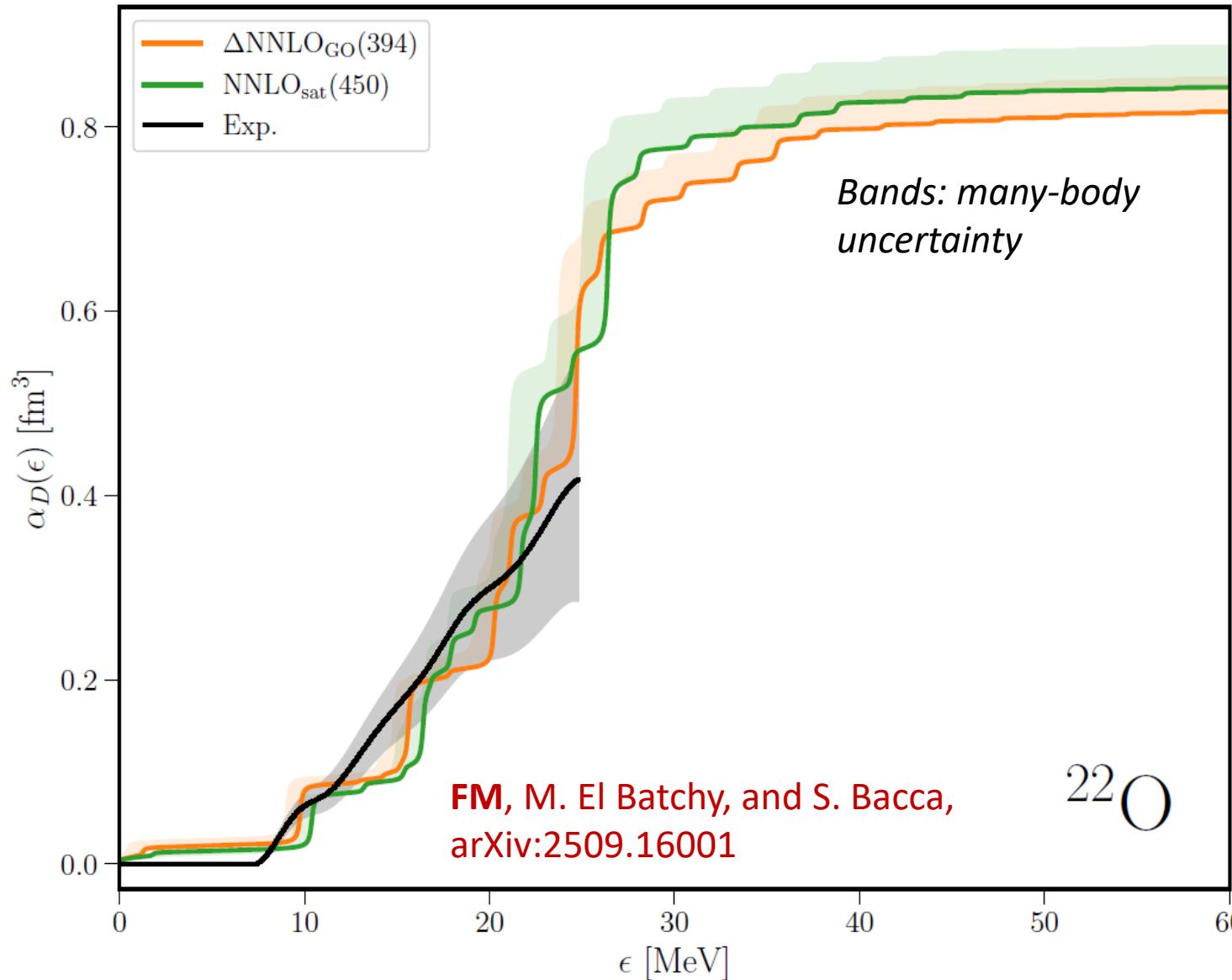


Calculations at the CCSDT-1 level

LIT-CC consistently predicts
low-lying strength near the
neutron threshold

The ^{22}O electric dipole polarizability

Polarizability running sum



$$\alpha_D(\epsilon) = 2\alpha\hbar c \int_0^\epsilon d\sigma \frac{L(\sigma, \Gamma')}{\sigma}$$

$$\Gamma' = 10^{-4} \text{ MeV}$$

Good agreement between theory and experiment (up to 25 MeV)

From the LIT to the response function

Problem: Determine the response function $R(\omega)$ from the LIT

$$L(\sigma, \Gamma) \quad \xrightarrow{\text{red arrow}} \quad R(\omega) \quad \text{III-posed numerical inversion!}$$

From the LIT to the response function

Problem: Determine the response function $R(\omega)$ from the LIT

$$L(\sigma, \Gamma) \quad \xrightarrow{\text{red arrow}} \quad R(\omega) \quad \text{Ill-posed numerical inversion!}$$

Ansatz: Basis expansion with hyperparameters α, β, ν

$$\tilde{R}(\omega) = \sum_{n=1}^{\nu} c_n f_n(\omega; \alpha, \beta) \quad f_n(\omega; \alpha, \beta) = \omega^{\alpha} e^{-\frac{\omega}{n\beta}}$$

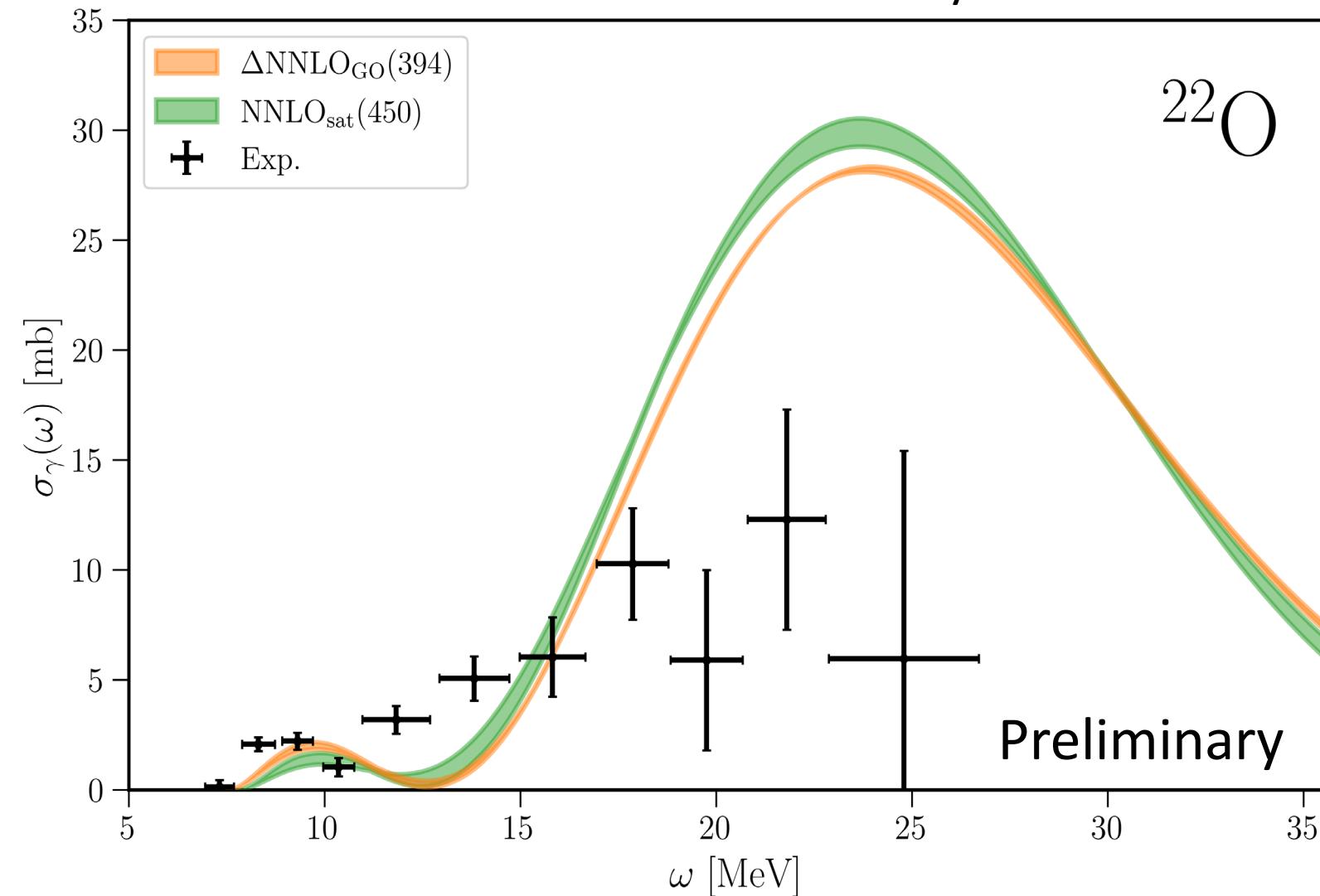
$$\chi^2 \text{ minimization} \quad \chi^2 = |L(\sigma, \Gamma) - \tilde{L}(\sigma, \Gamma)|^2$$

Uncertainties on $R(\omega)$ from varying the number of basis states ν and width Γ



^{22}O cross sections

Credit: Miriam El Batchy



$$\sigma_\gamma(\omega) = 4\pi^2 \alpha \omega R(\omega)$$

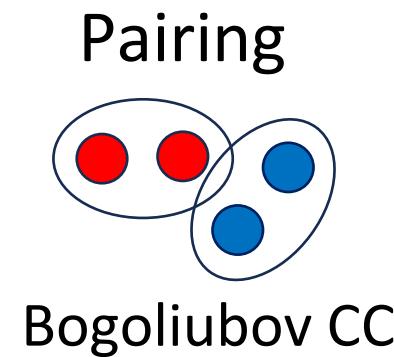
Exp. photoneutron cross sections from [[Phys. Rev. Lett. 86, 5442 \(2001\)](#)]

Cross sections with quantified **uncertainties** from the **inversion** of the LIT

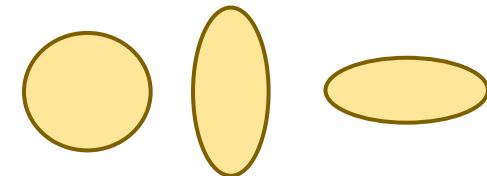
Structure and electric dipole polarizability of open-shell nuclei

Towards open-shell nuclei

Symmetry-breaking approaches



Deformation



Deformed CC

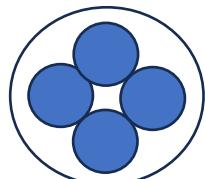
Tichai et al., Phys. Lett. B **851**, 138571 (2024)
Hagen et al., Phys. Rev. C **105**, 064311 (2022)

Towards open-shell nuclei

Symmetry-breaking approaches

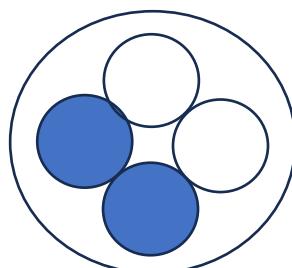
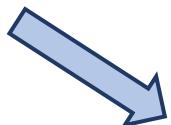
Equations-of-motion

Closed-shell

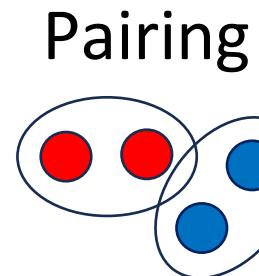


$$|\Psi_0\rangle = e^T |\Phi_0\rangle$$

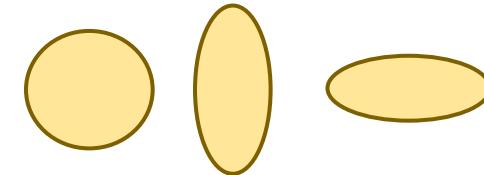
Two-particle-removed
(2PR) nucleus



Pairing
Bogoliubov CC



Deformation
Deformed CC



Tichai et al., Phys. Lett. B 851, 138571 (2024)
Hagen et al., Phys. Rev. C 105, 064311 (2022)

$$|\Psi_f^{(A-2)}\rangle = R_f^{(A-2)} |\Psi_0\rangle$$

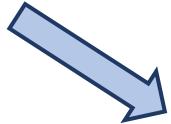
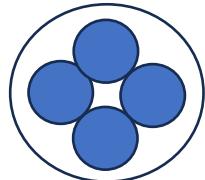
Towards open-shell nuclei

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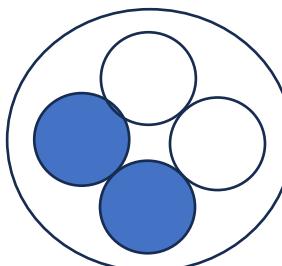
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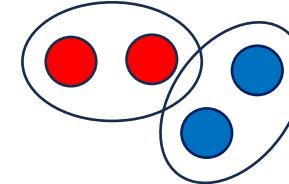
Two-particle-removed
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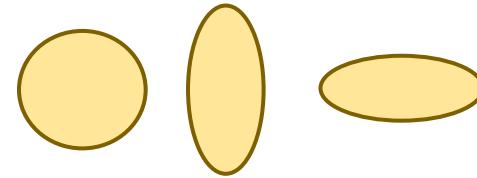
$$|\Psi_0\rangle = e^T |\Phi_0\rangle$$



Pairing
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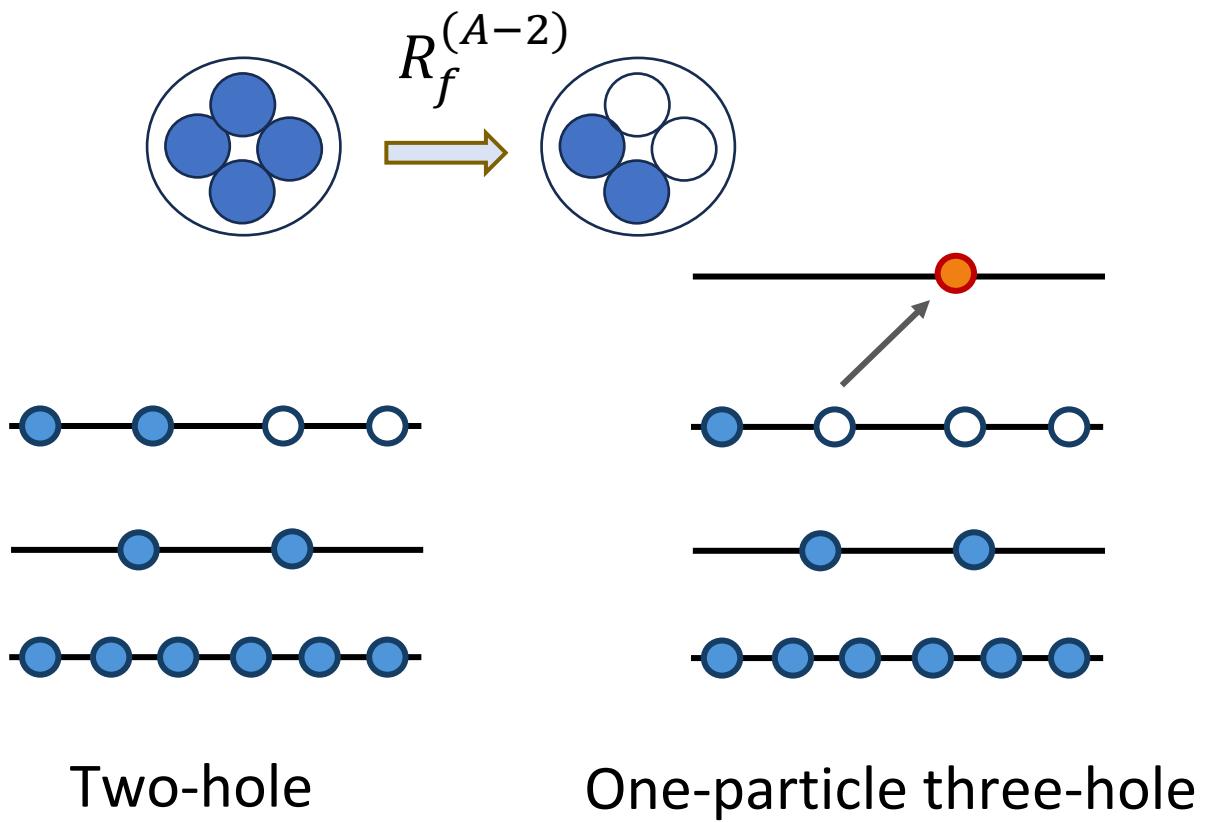


Tichai et al., Phys. Lett. B 851, 138571 (2024)
Hagen et al., Phys. Rev. C 105, 064311 (2022)

Excitation operator

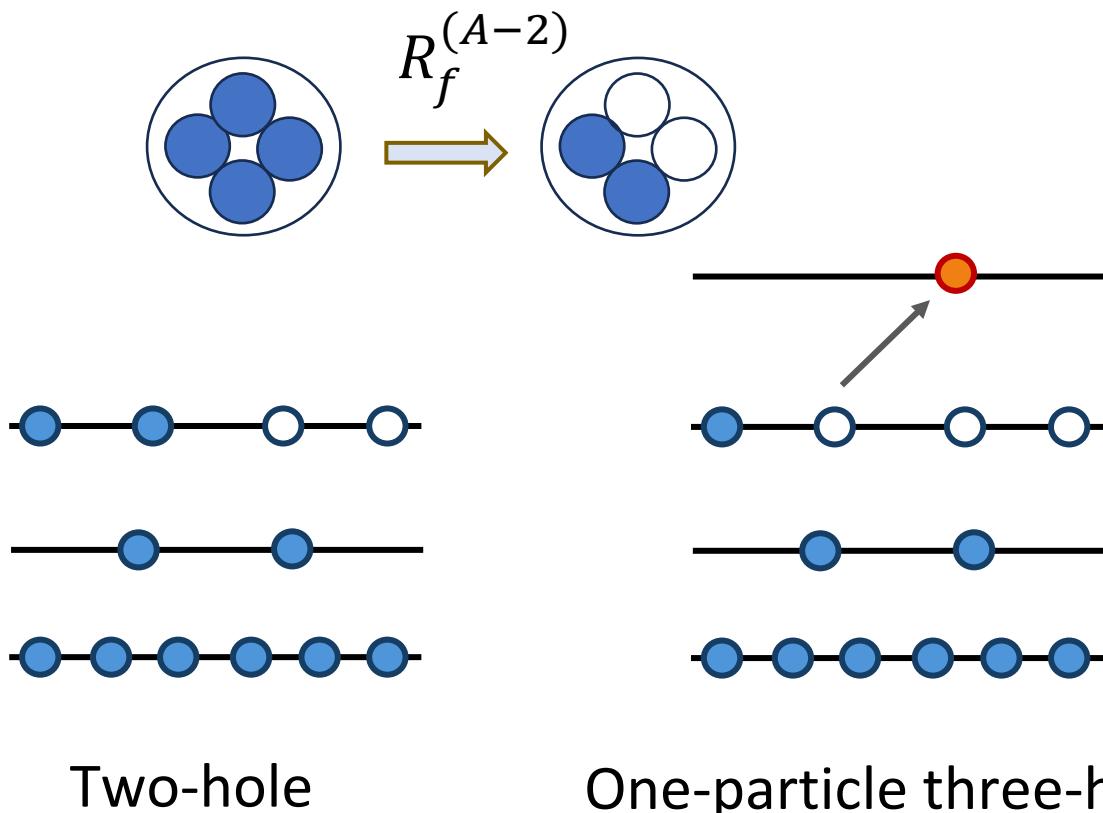
$$|\Psi_f^{(A-2)}\rangle = R_f^{(A-2)} |\Psi_0\rangle$$

The two-particle-removed ansatz



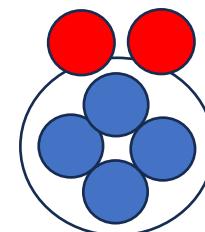
$$R_f^{(A-2)} = \frac{1}{2} \sum_{ij} r_{ij} c_j c_i + \frac{1}{6} \sum_{ijk a} r_{ijk}^a c_a^\dagger c_k c_j c_i$$

The two-particle-removed ansatz



$$R_f^{(A-2)} = \frac{1}{2} \sum_{ij} r_{ij} c_j c_i + \frac{1}{6} \sum_{ijk a} r_{ijk}^a c_a^\dagger c_k c_j c_i$$

See also

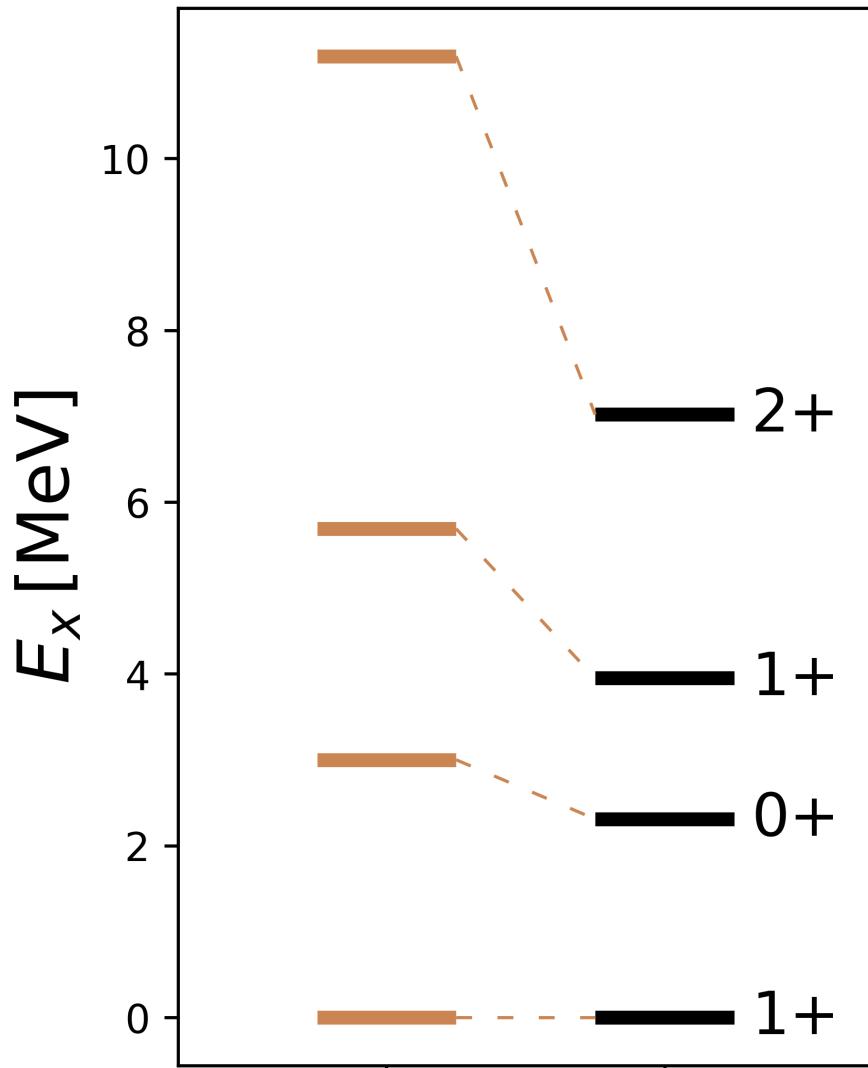


Two-particle-attached (2PA)

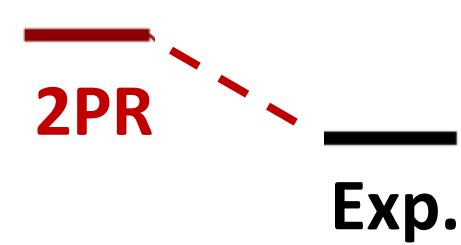
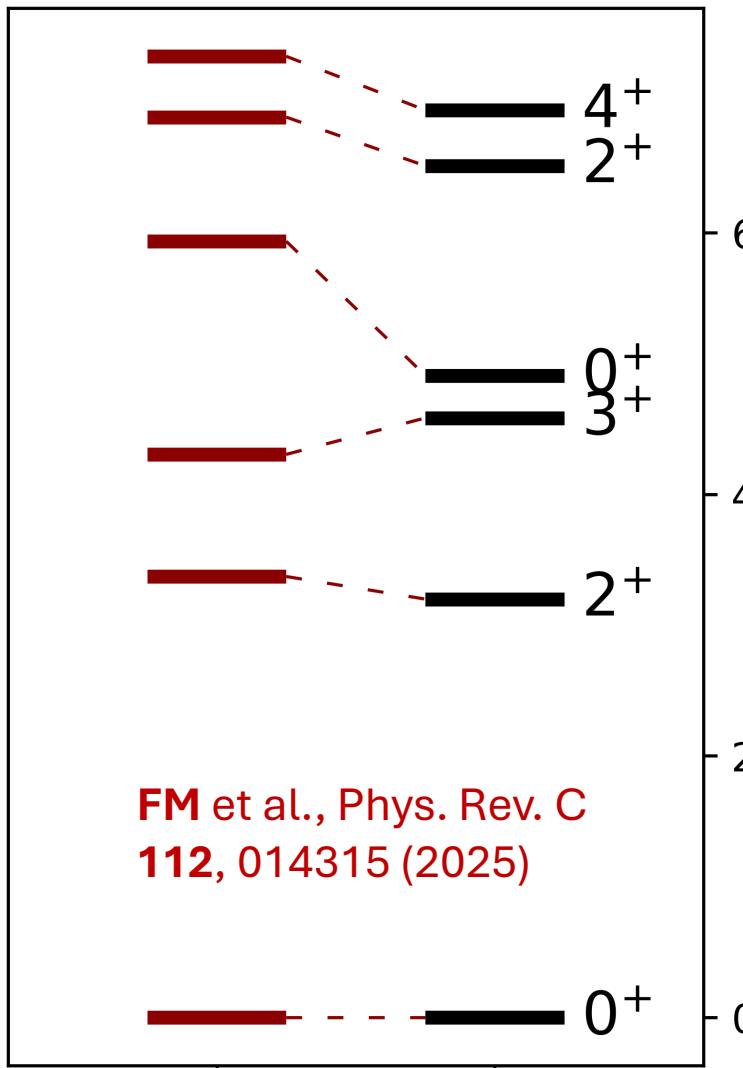
Jansen, Phys. Rev. C **88**, 024305 (2013)
Bonaiti et al., PRC **110**, 044306 (2024)

Low-lying spectra

^{14}N

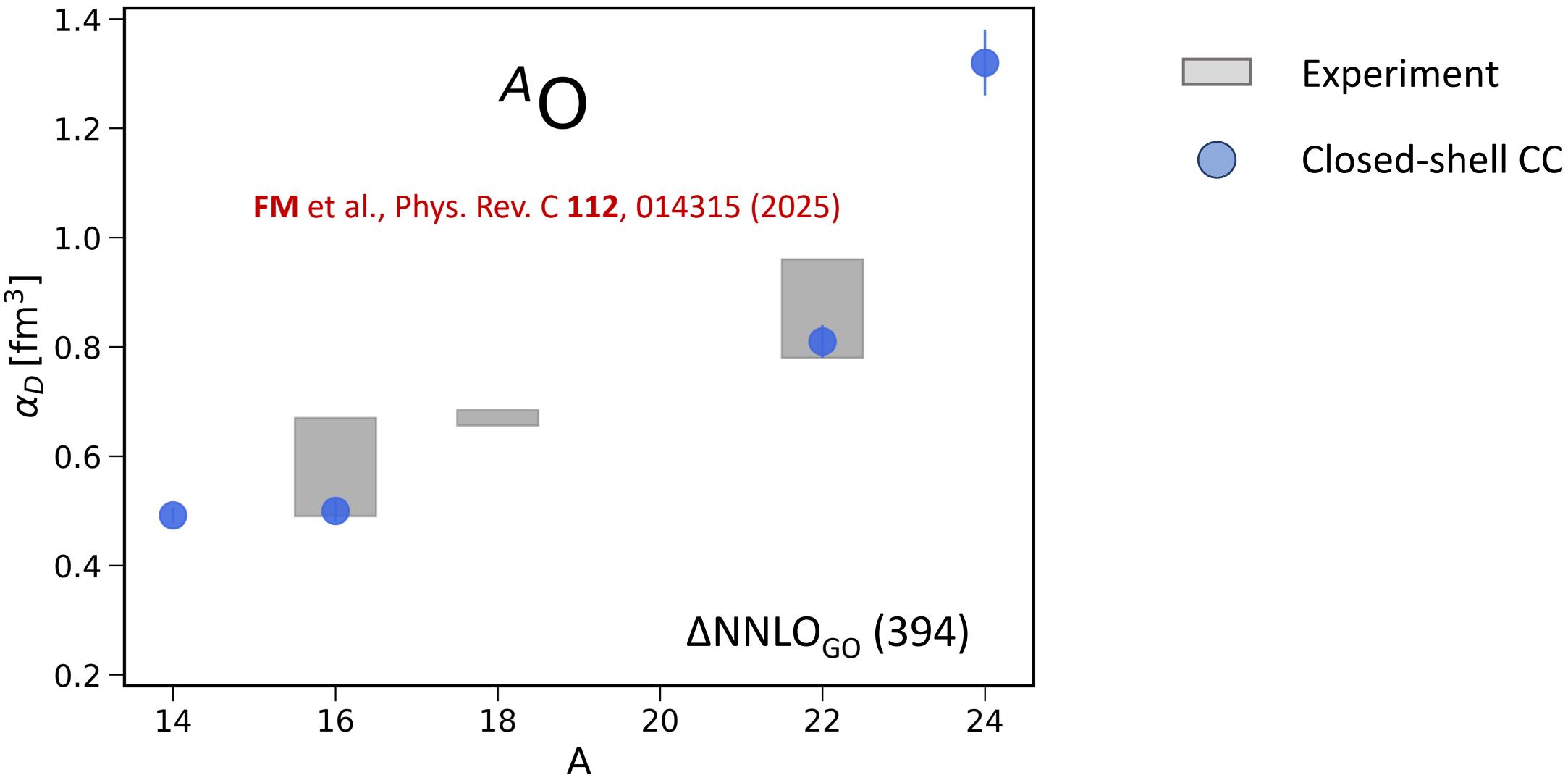


^{22}O

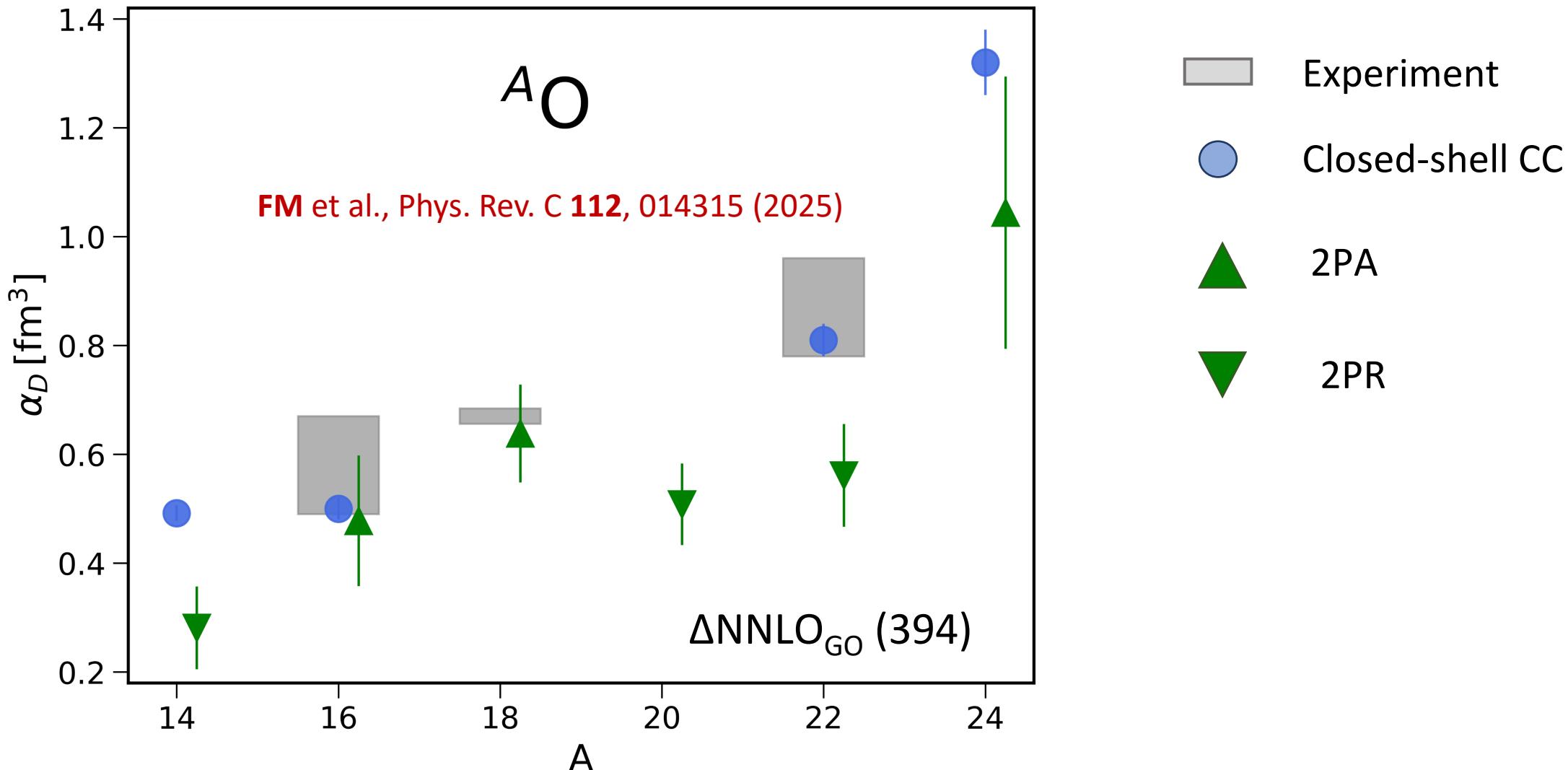


2PR with the $\Delta\text{NNLO}_{\text{GO}} (394)$ interaction describes well several **excited states**

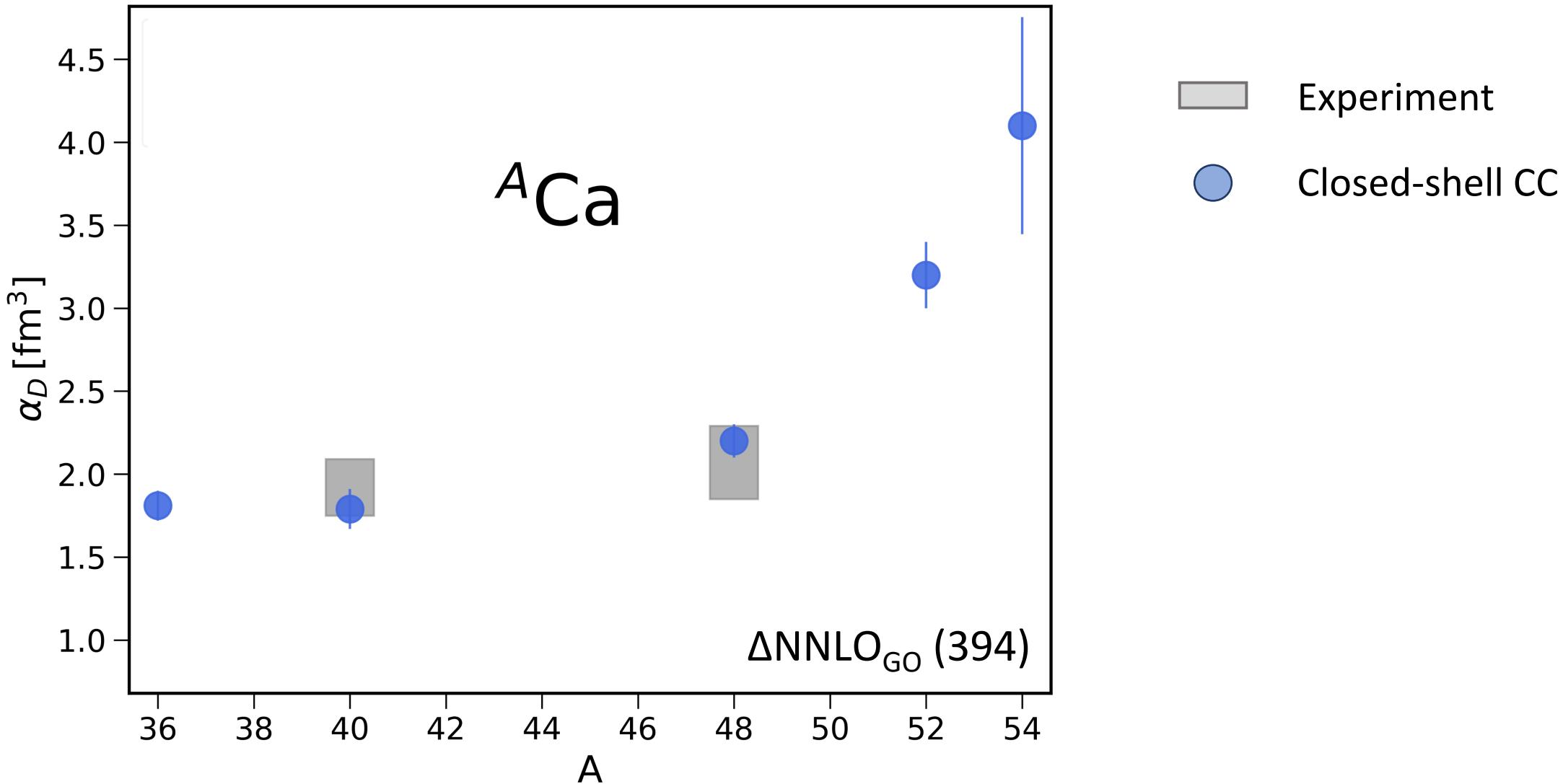
Electric dipole polarizability in O isotopes



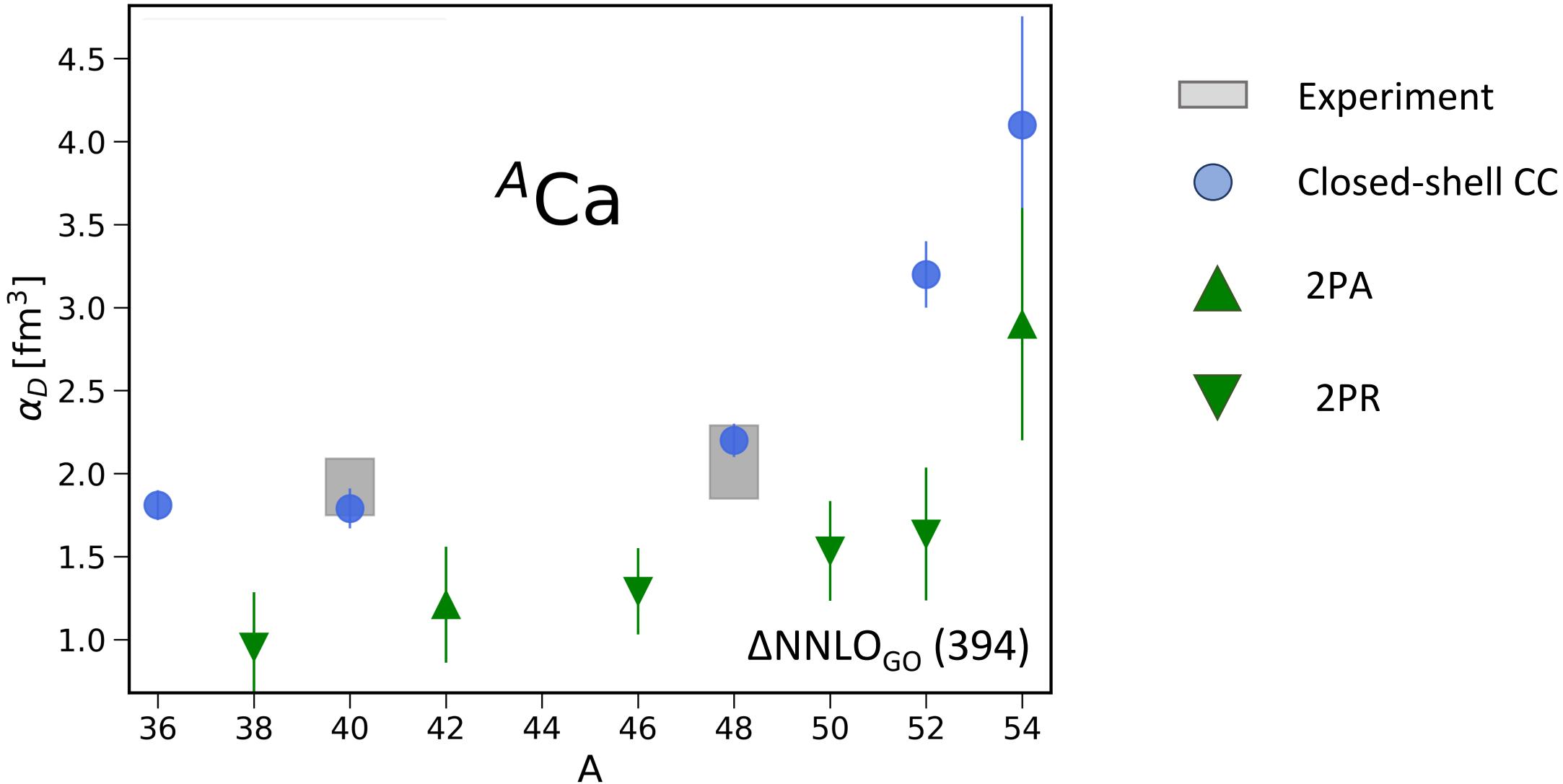
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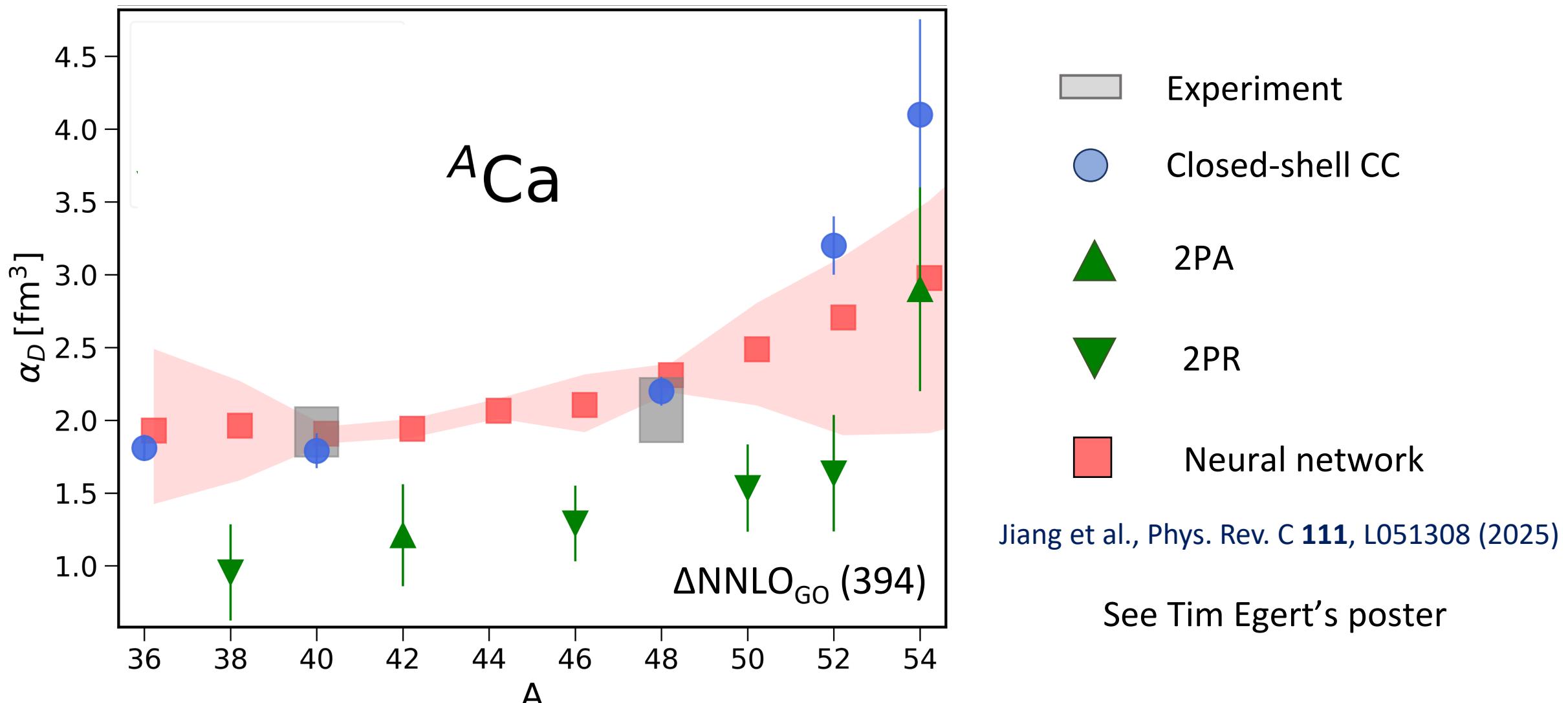
Electric dipole polarizability in Ca isotopes



Electric dipole polarizability in Ca isotopes



Electric dipole polarizability in Ca isotopes



Conclusions and perspectives

- We have applied coupled-cluster to the dipole response and electric dipole polarizability of the neutron-rich O22 isotope
- We have extended the method to the ground state, excited states and electric dipole polarizability of open-shell nuclei close to magicity
- Future developments: response of odd nuclei, response functions with theory uncertainties ...

Stay tuned!

Many-body Theory: Nuclear Physics Meets Quantum Chemistry

Aug 24th – Sep 4th 2026



Mainz Institute for
Theoretical Physics

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